

# High-order Spectral Difference Method for Simulating Unsteady Flow past a plunging airfoil

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Vortex dominated flow workshop  
National Institute of Aerospace  
Virginia, June, 2009

# OVERVIEW

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- **Introduction**
- **Spectral Difference method**
- **Validation**
- **Vortex Shedding**
- **Subsonic Flow past a plunging airfoil**

# Higher order Methods (DG,SV,SD)



- Low numerical dissipation
- Less numerical dispersion
- Unstructured grids
- CPU efficient/easy to parallelize
- Not memory intensive
- Easy to program, universal construction by placing unknown points in a geometrical similar manner
- SD attains a simpler form and higher efficiency than DG and SV.
- Flexible (grid-independence, hp adaptation, moving boundary, deformable grid)

# History of Spectral Difference Method



## Two early papers on the SD method

- D. A. Kopriva and J. H. Kollias, A conservative staggered-grid chebyshev multidomain method for compressible flows. J. Comput. Phys. 125 (1996), p. 244. **STRUCTURED STAGGERED GRID**
- Yen Liu , Marcel Vinokur , Z. J. Wang, Spectral difference method for unstructured grids I: basic formulation, Journal of Computational Physics, v.216 p.780-801,2006. **UNSTRUCTURED MULTIVARIATE FORMULATION**

## Our papers on the SD method

- Wang, Z.J., Liu, Y., May, G., Jameson, A.: "Spectral Difference Method for Unstructured Grids II: Extension to the Euler Equations" , J. Sci. Comput. 32 (1) pp. 54-71, July, 2007. **EXTENSION TO EULER EQUATIONS**
- C. Liang, S. Premasathan, A. Jameson, "High-order accurate simulation of flow past two side-by-side cylinders with Spectral Difference method", 2009, vol 87, pp. 812-817, Journal of Computers and Structures. **EXTENSION TO 2D Viscous flow on quadrilateral elements**
- C. Liang, A. Jameson and Z. J. Wang, "Spectral Difference method for two-dimensional compressible flow on unstructured grids with mixed elements", Journal of Computational Physics, vol 228, pp 2847-2858, 2009. **EXTENSION TO 2D Viscous flow on mixed elements**

# Development of the spectral difference method in Stanford University



- 2D quadrilateral/triangular element
- 3D hexahedral elements
- 4<sup>th</sup>-order and higher on unstructured grids
- Implicit LU-SGS time stepping
- P-multigrid method
- Moving and deforming grids
- Artificial viscosity for shock capturing
- Parallelization (MeTis/MPI) of 3D code

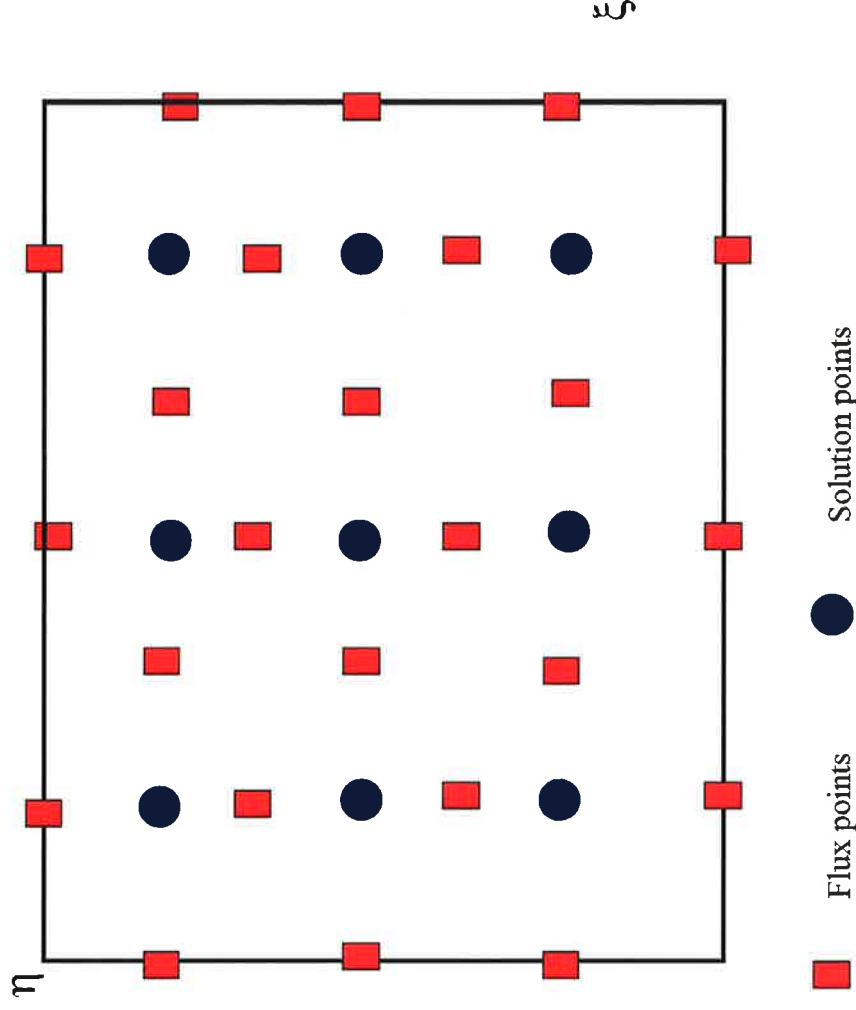


# High-order spectral difference method



- The figure shows one way to locally construct a third-order spectral difference scheme for a grid cell.
- 9 solution points are used.
- 24 flux points are employed.
- The reconstructed field using polynomials is continuous within the cell but discontinuous across the cell boundaries

$$\frac{dq}{dt} + df(q)/dx = 0$$



- Flux points store  $f(q)$
- Solution points store  $q$

# Choosing solution and flux points



- N solution points (one can choose arbitrarily)

$$(a) \quad X_s = \frac{1}{2} \left[ 1 - \cos \left( \frac{2s-1}{2N} \cdot \pi \right) \right], s = 1, 2, \dots, N. \quad \text{Gauss points}$$

- N+1 flux points

$$(b) \quad X_{s+1/2} = \frac{1}{2} \left[ 1 - \cos \left( \frac{s}{N} \cdot \pi \right) \right], s = 0, 1, \dots, N. \quad \begin{array}{l} \text{Chebyshev-Gauss-} \\ \text{Lobatto points} \end{array}$$

Or

$$(c) \quad \text{Roots of the equation: } P_n(\xi) = 0 \quad \begin{array}{l} \text{Legendre-Gauss} \\ \text{quadrature points and two} \\ \text{end points of 0 and 1.} \end{array}$$

$$P_{-1}(\xi) = 0$$

$$P_0(\xi) = 1 \quad P_n(\xi) = \frac{2n-1}{n} (2\xi-1) P_{n-1}(\xi) - \frac{n-1}{n} P_{n-2}(\xi)$$



## Stability of the Spectral Difference method

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- K Abeele, C Lacor and Z. J. Wang (2008), *On the stability and accuracy of the spectral difference method*, Journal of Scientific Computing, vol 37, pp. 162-188.
- Antony Jameson, (2009), *A proof of the stability of the spectral difference method for all orders of accuracy*, Aerospace Computing Lab report, ACL-2009-1, Stanford University, March 2009.
- Both confirm stability when the interior flux points are at the zeros of the Legendre polynomial



# Solution and Flux reconstruction



Degree N-1 polynomial through solution points using Lagrange basis polynomial

$$h_i(X) = \prod_{s=1, s \neq i}^N \left( \frac{X - X_s}{X_i - X_s} \right)$$

Degree N polynomial through flux points using Lagrange basis polynomial

$$l_{i+1/2}(X) = \prod_{s=0, s \neq i}^N \left( \frac{X - X_{s+1/2}}{X_{i+1/2} - X_{s+1/2}} \right)$$

Reconstructed solution as well as flux polynomials within an element are written as tensor products of three 1D polynomials

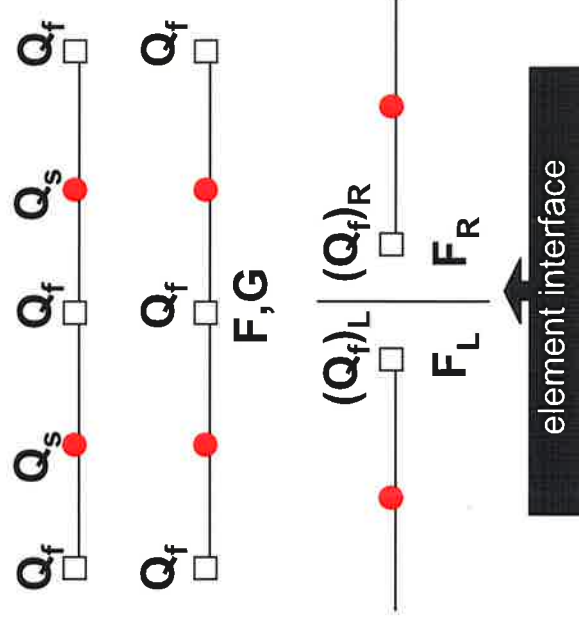
$$Q(\xi, \eta, \beta) = \sum_{k=1}^N \sum_{j=1}^N \sum_{i=1}^N \frac{\tilde{Q}_{i,j,k}}{|J_{i,j,k}|} h_i(\xi) \cdot h_j(\eta) \cdot h_k(\beta)$$

$$\tilde{F}(\xi, \eta, \beta) = \sum_{k=1}^N \sum_{j=1}^N \sum_{i=0}^N \tilde{F}_{i+1/2,j,k} l_{i+1/2}(\xi) \cdot h_j(\eta) \cdot h_k(\beta)$$

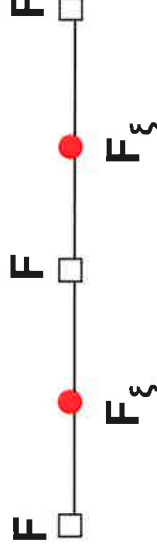
# Calculation of inviscid flux derivatives



- Given the conservative variables at the solution points, the conservative variables are extrapolated to the flux points
- The inviscid fluxes at the interior flux points are computed
- The inviscid fluxes at the element interfaces are computed using an approximate Riemann solver
- The derivative of the fluxes are computed at the solution points according to



$$\left( \frac{\partial \tilde{F}}{\partial \xi} \right)_{i,j} = \sum_{r=0}^N \tilde{F}_{r+1/2,j} \cdot l'_{r+1/2}(\xi_i)$$

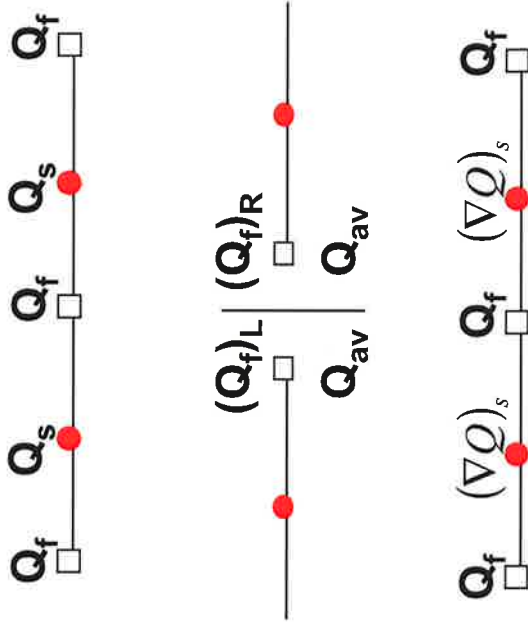


# Calculation of viscous flux derivatives (1)



- The viscous fluxes depend on both the solution and the solution gradient at the flux points
- Reconstruct conservative variables at flux points using values at solution points
- At flux points on element interfaces, compute the average of left and right solution values
- Calculate the gradient of  $Q$  at the solution points from  $Q$  at flux points

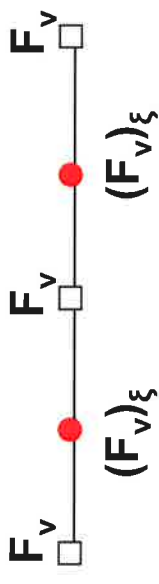
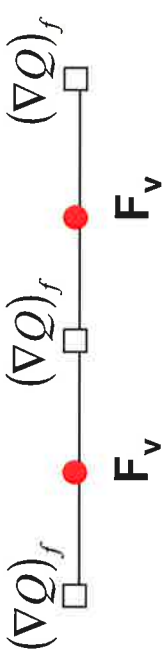
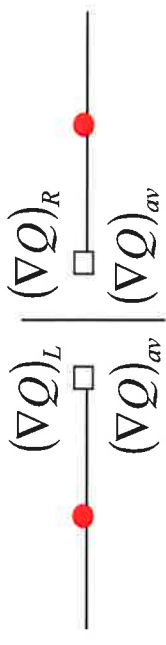
$$F_v = F_v(Q_f, \nabla Q_f)$$



# Calculation of viscous flux derivatives (2)



- Reconstruct  $\nabla Q$  from solution points to flux points
- At flux points on element interfaces, compute the average of left and right  $\nabla Q$  values
- Compute viscous fluxes at flux points
- Compute viscous flux derivatives at solution points
- We are ready to compute the NS equation, using a five-stage R-K scheme



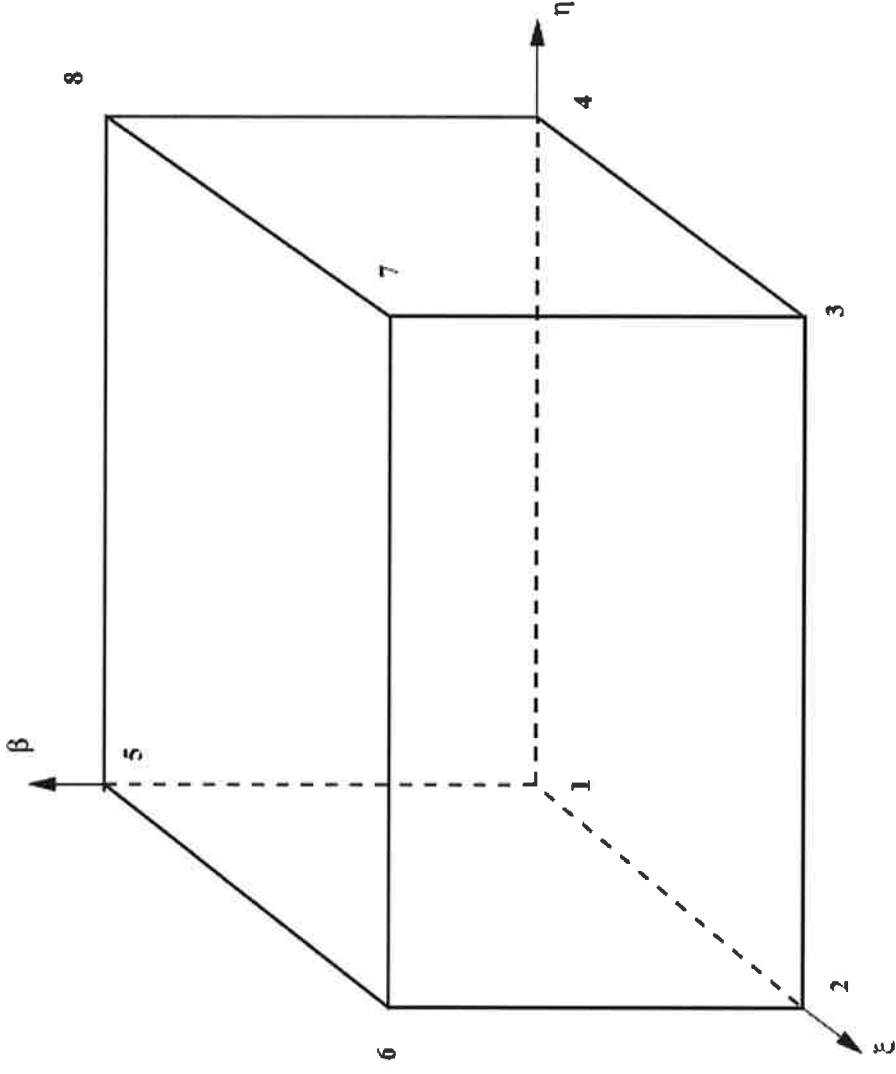
$$\frac{\partial Q}{\partial t} + \nabla F_e(Q) - \nabla F_v(Q, \nabla Q) = 0$$

# Hexahedral grid mapping



$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \sum_{i=1}^K M_i(\xi, \eta, \beta) \begin{pmatrix} x_i \\ y_i \\ z_i \end{pmatrix}$$

□ Mapping from physical hexahedral grid elements to a standard computational element

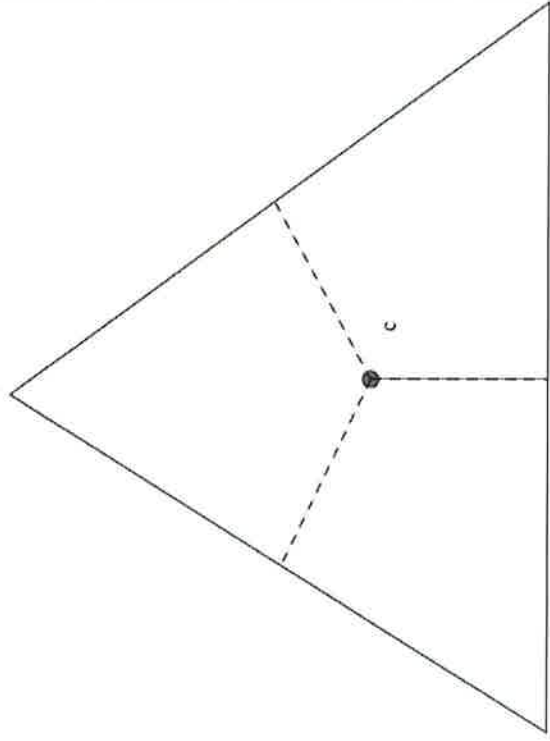


$$(0 \leq \xi \leq 1, 0 \leq \eta \leq 1 \text{ and } 0 \leq \beta \leq 1)$$

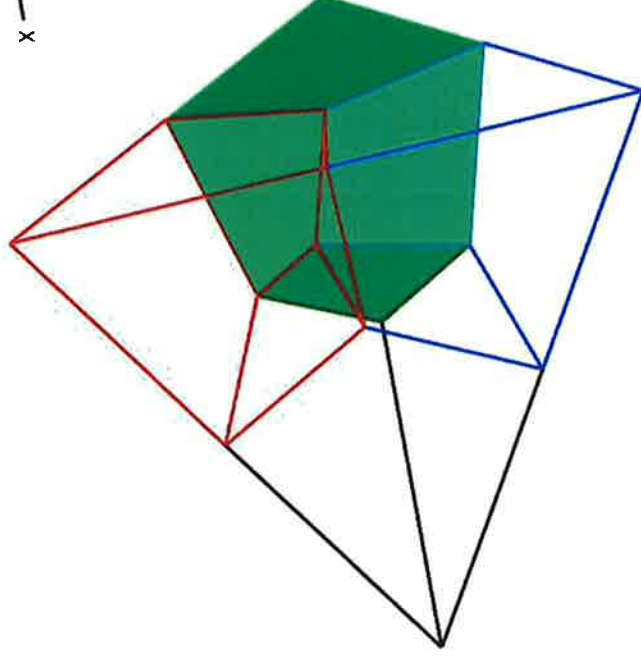
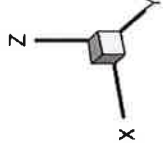




## Splitting triangulated grids to quadrilateral grids in 2D and 3D

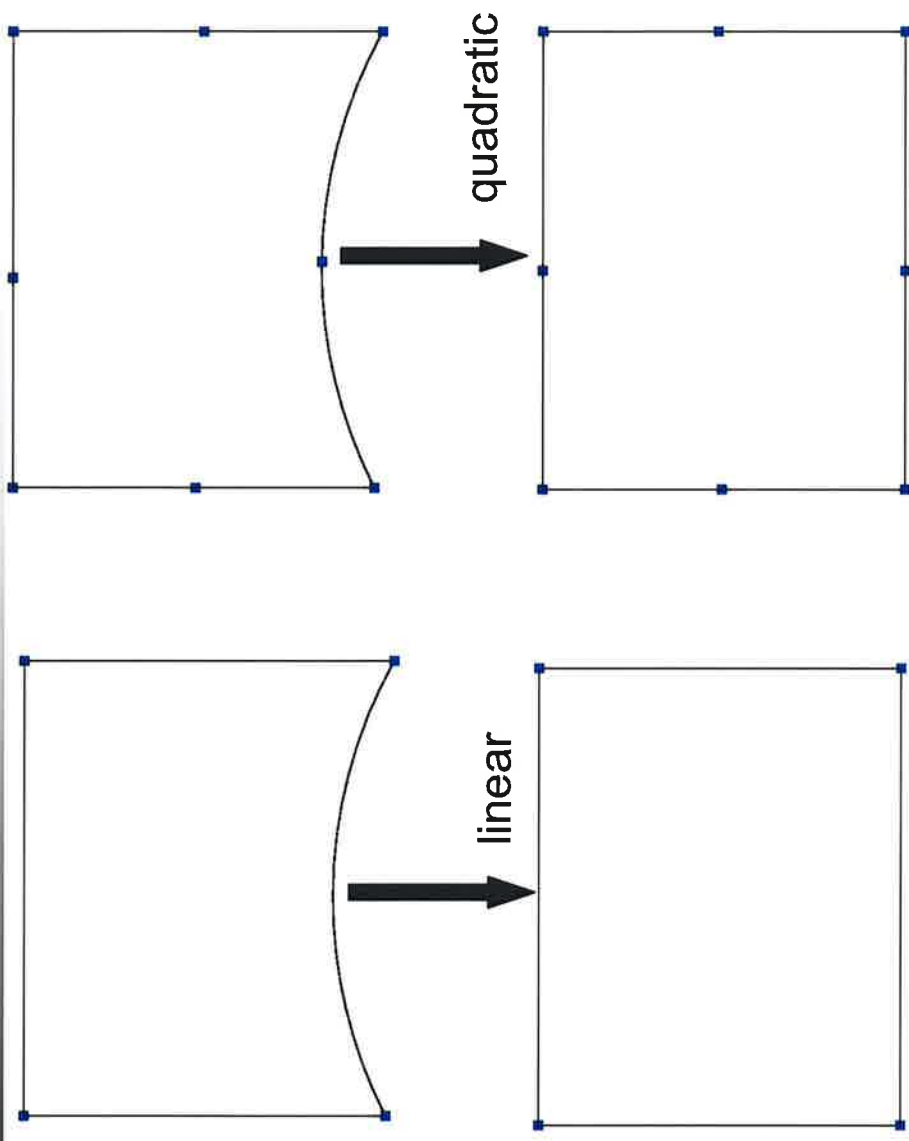


(a) A triangle split into three quadrilaterals



(b) A tetrahedron split into four hexahedra

# Curved wall representation



$$\begin{pmatrix} x \\ y \end{pmatrix} = \sum_{i=1}^K M_i(\xi, \eta) \begin{pmatrix} x_i \\ y_i \end{pmatrix}$$

M is shape function; K = 8, 20 for 3D linear, quadratic mapping

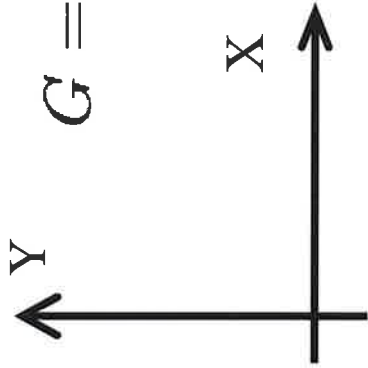
K=4,8 for 2D linear, quadratic mapping

# Extension to Moving and Deformable Grids (I)



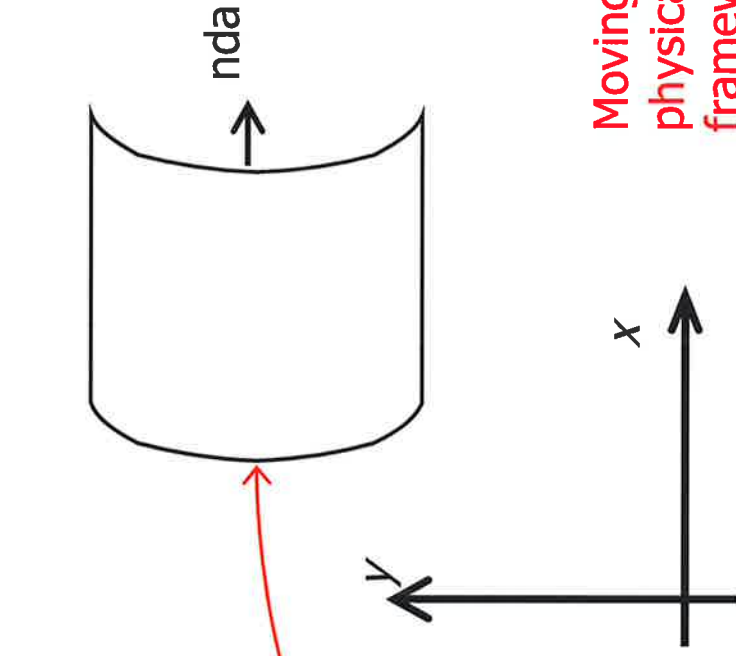
$$(x, y, z) = M(X, Y, Z, t)$$

$$G = \nabla_X M = \frac{\partial(x, y, z)}{\partial(X, Y, Z)}$$



$\rightarrow NdA$

$G, g, w$



$\rightarrow nda$

Moving physical framework

Reference framework

# Extension to Moving and Deformable Grids (II)



Define a mapping velocity  
:

$$\frac{\partial M(X, Y, Z, t)}{\partial t} = w$$

$$\frac{\partial \hat{Q}}{\partial t} |_X + \nabla_X \hat{F}_e(\hat{Q}) - \nabla \hat{F}_v(\hat{Q}, \nabla \hat{Q}) = 0$$

$$g = \det(G) \quad \hat{Q} = gQ$$

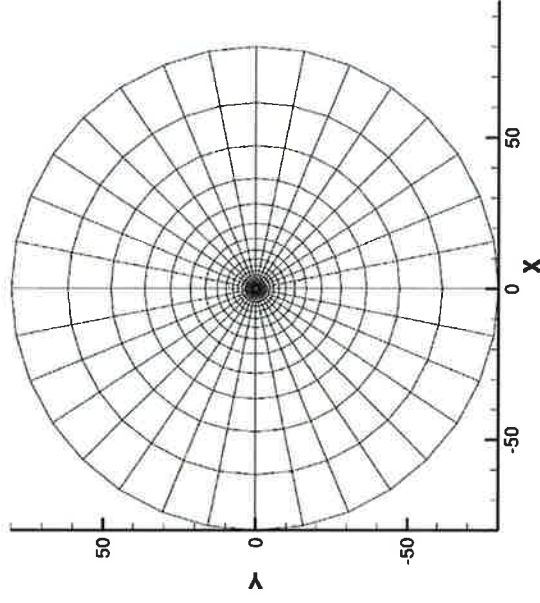
$$\hat{F} = gG^{-1}F - \hat{Q}G^{-1}w$$



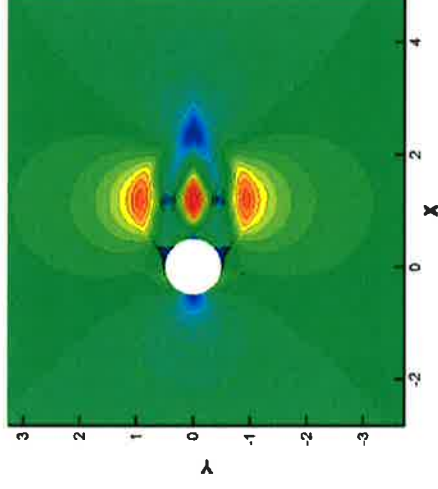
# VALIDATION



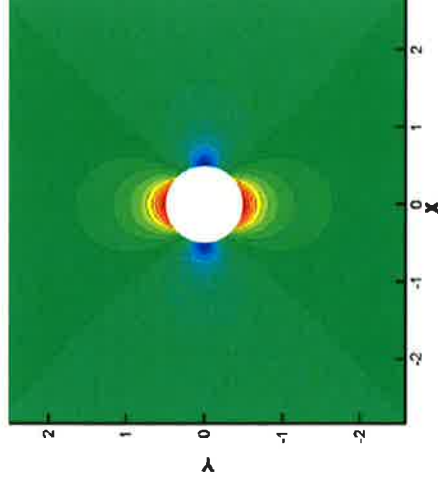
# Mach contour around a circle



4<sup>th</sup>-order SD method on a grid  
with 32x32 cells



(a) linear wall boundary



(b) cubic wall boundary

- Free stream Mach = 0.2
- Inviscid Euler solver

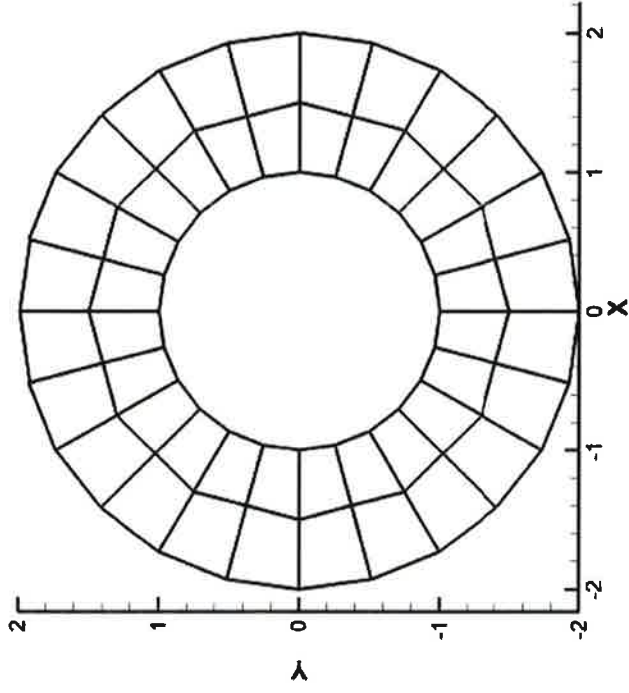
# Effect of interface flux on Cd



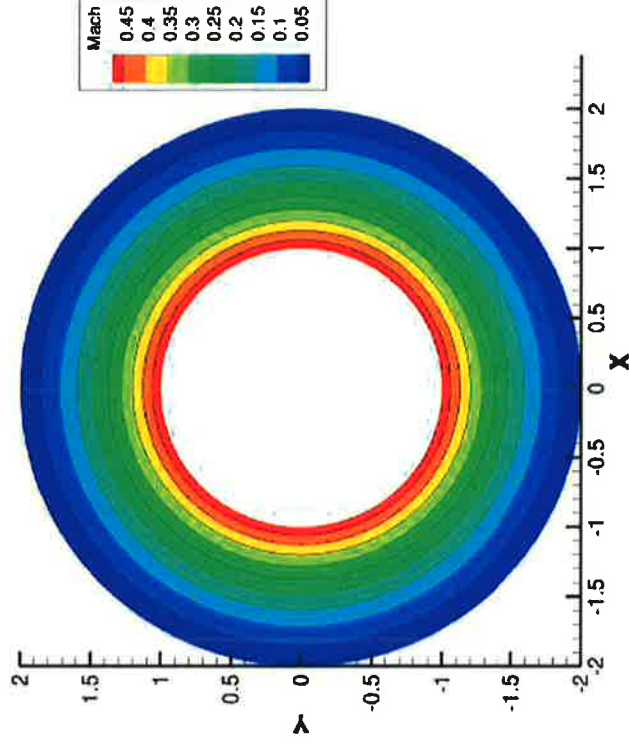
Riemann	ordercell no.	Wall	Mach	$\Delta_t u_\infty / D$	$C_d$
CUSP	4th	640	quadratic	0.2	2e-4 -1.86e-5
AUSM	4th	640	quadratic	0.2	2e-4 -4.39e-5
Roe	4th	640	quadratic	0.2	2e-4 -1.03e-5
flux vector split	4th	640	quadratic	0.2	2e-4 -1.18e-5
Scalar diffusion	4th	640	quadratic	0.2	2e-4 -8.8e-6

□ Inviscid flow past a circle

# Compressible Taylor-Couette Flow



(a) Grid



(b) Mach contour

Mach = 0.5,  $Re=10$ , isothermal for inner cylinder and adiabatic wall for outer cylinder

# Verification of order of accuracy




No. of elements	No. of DOFs	L2-error	Order
3rd order SD			
48	432	8.896E-04	-
192	1728	1.002E-04	3.15
768	6912	1.084E-05	3.21
4th order SD			
48	768	1.4815E-04	-
192	3072	1.0036E-05	3.88
768	12288	6.5746E-07	3.93

Exact solution for  
angular velocity

$$v_{\theta} = \Omega_i r_i \frac{r_o}{r_i} - \frac{r}{r_o} \frac{r_i}{r_o}$$

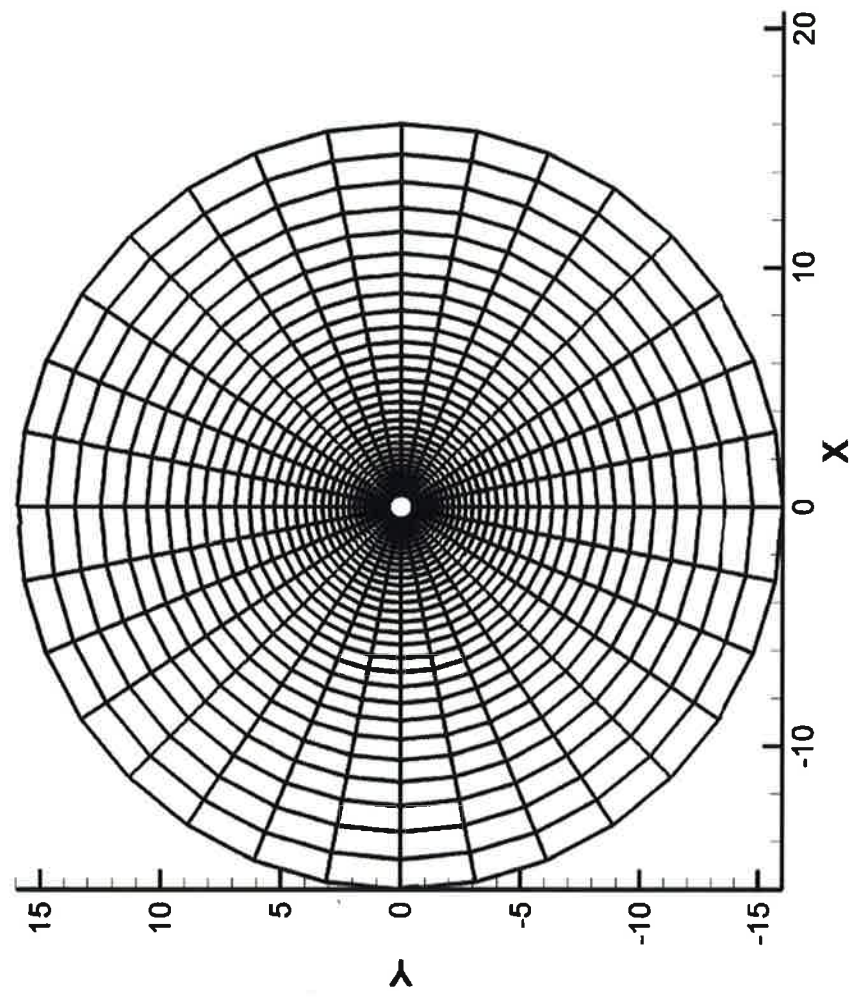


# Studies of Vortex Shedding Past a Plunging Cylinder

 E. GUILMINEAU and P. QUEUTEY,  
*A NUMERICAL SIMULATION OF VORTEX SHEDDING  
FROM AN OSCILLATING CIRCULAR CYLINDER*  
Journal of Fluids and Structures  
Vol 16, 2002, pp. 773-794



# Computational grid for oscillating cylinder



- Our simulation uses only 32x32 grid
- The third-order SD method
- Total degrees-of-freedom = 96x96
- Quadratic curved wall



## vorticity shedding behind an oscillating cylinder

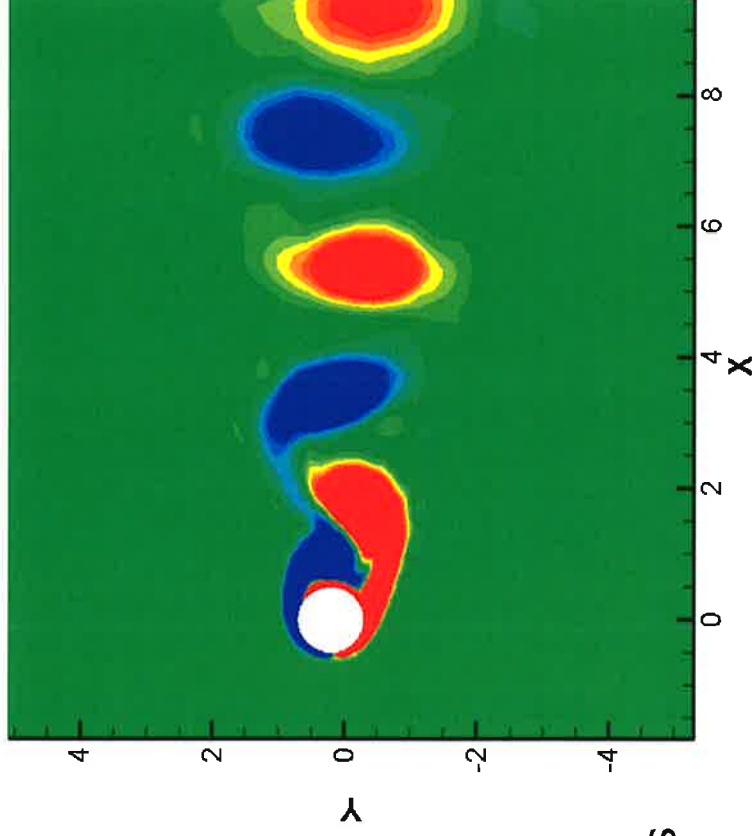
- $Y(t) = Ae \cos(2\pi fe t)$   
for cylinder transverse motion

- $Ae = 0.2D$

- $fe = 1.1 \text{ fn}$

- $-1 < \omega d/U < 1$

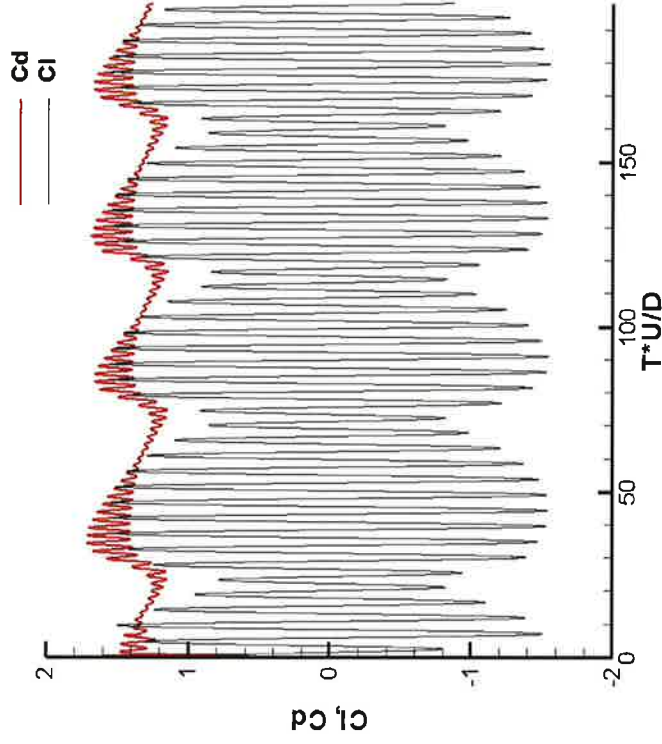
- $Re = 185$



- The cylinder motion is identical to the one reported in E. GUILMINEAU and P. QUEUTEY, Journal of Fluids and Structures Vol 16, 2002, pp. 773-794

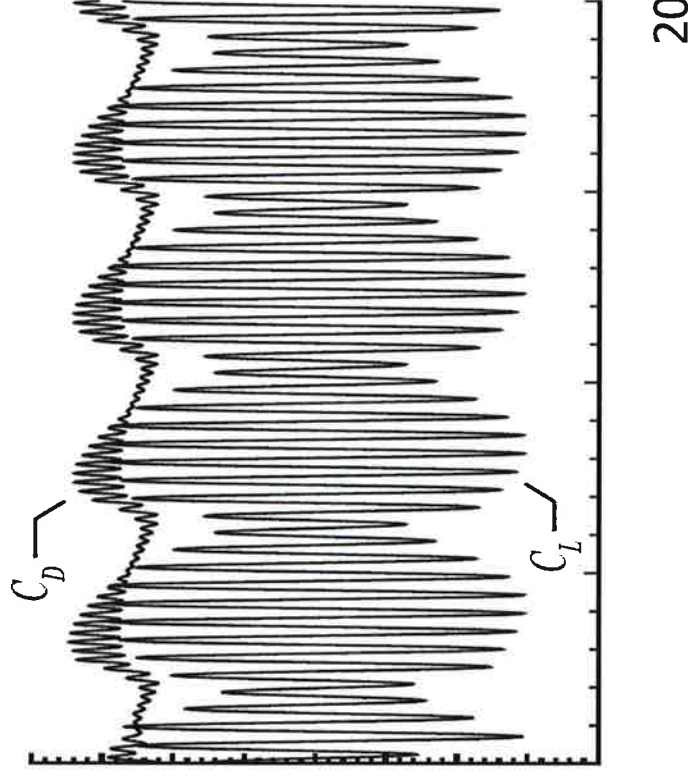
- We obtained nearly identical results using a grid with only 1024 cells compared to their 48,000 cells

# Fluid-exerted forces



(a) Third order spectral difference

Total cell : 32x32



(b) E. GUILMINEAU and P. QUEUTEY

Journal of Fluids and Structures

Vol 16, 2002, pp. 773-794

Total cell : 240x200

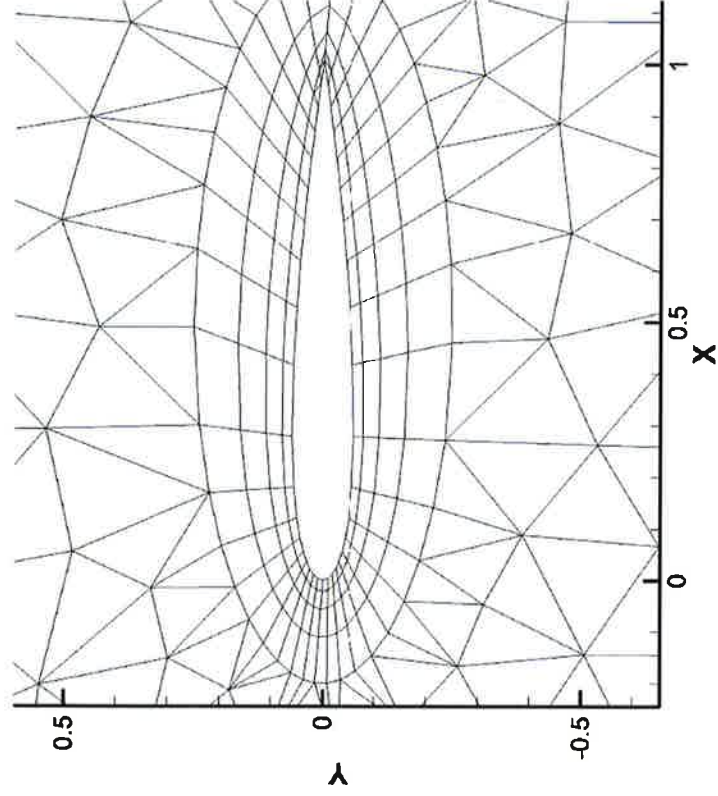
**The predicted force coefficients are in a good agreement with Guilmineau and Queutey (2002) and much less computational cells are used.**



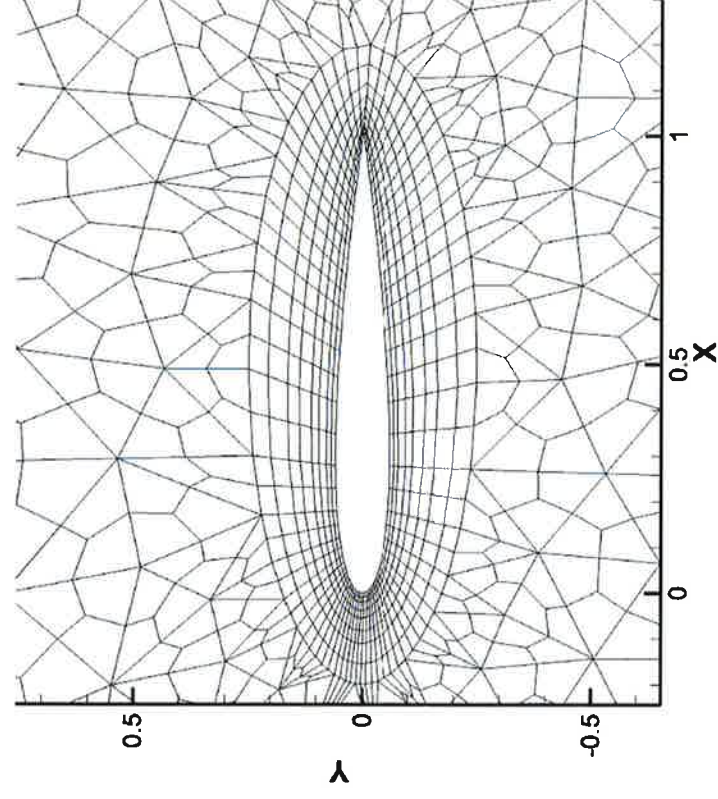
# Studies of plunging airfoils

□ K. D. Jones, C. M. Dohring and M. F. Platzer,  
Experimental and computational investigation of the  
Knoller-Betz effect,  
*AIAA Journal*, Vol 36, pp. 780-783.

# NACA 0012 grid with mixed elements



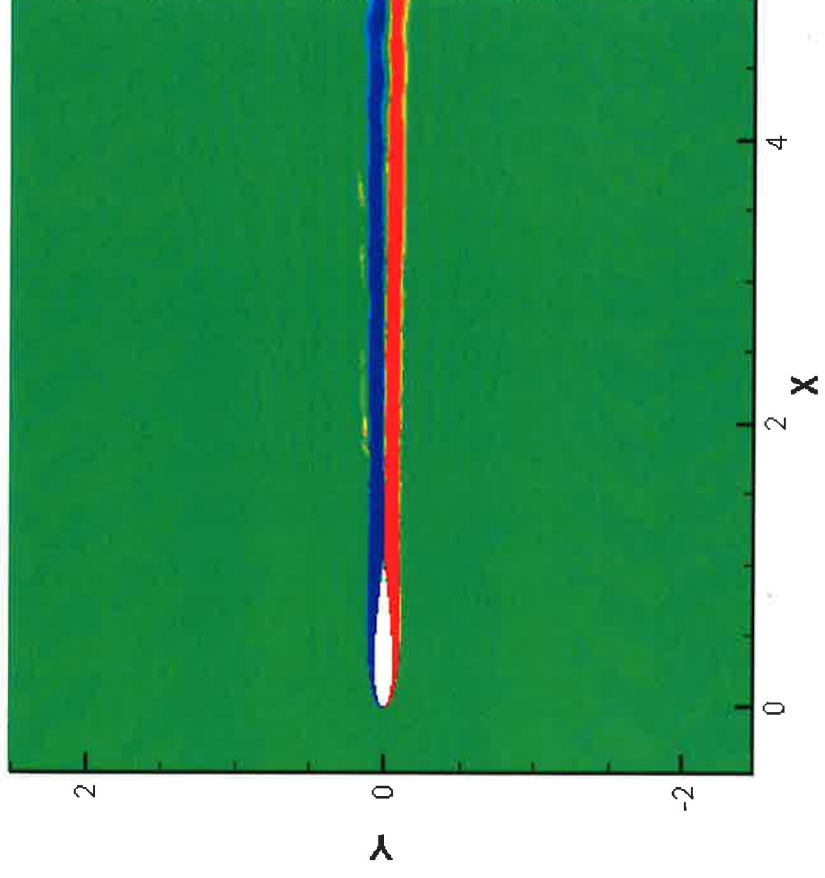
(a) Initial hybrid mesh: 'Mesh A'



(b) One-level h-refinement: 'Mesh B'



# Viscous flow past a stationary airfoil



- $Re = 1850$
- $Mach = 0.2$
- Vortices:

$$-6 \leq \omega c / U_\infty \leq 6$$

Steady flow solution  
is reached eventually

# Slow plunging airfoil case



■  $Y(t) = Ae \sin(\omega t)$

■  $\Omega = 1.15$

■  $Ae = 0.08 c$

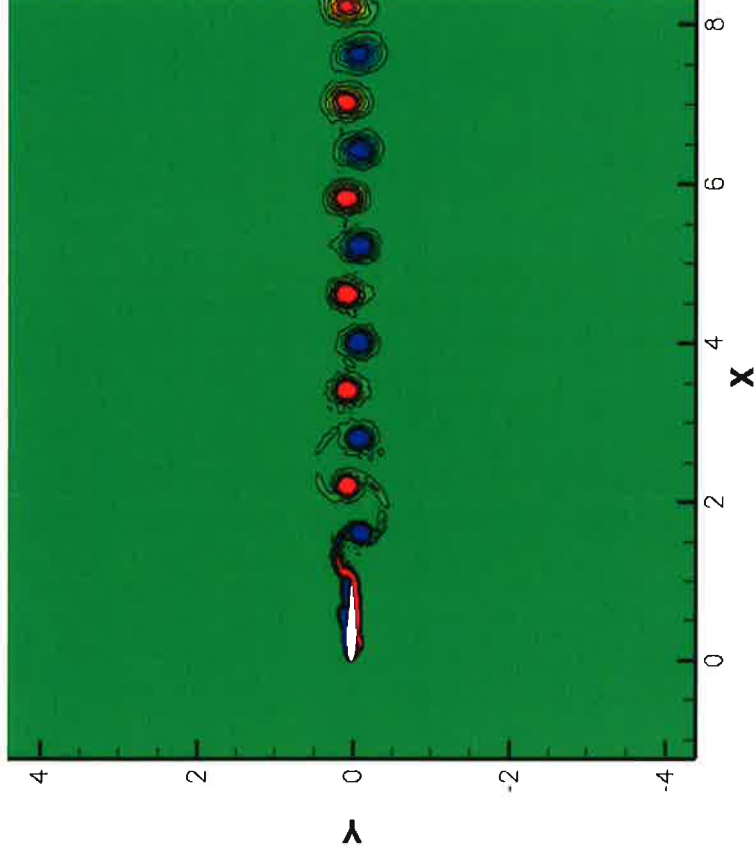
■ Chord  $c = 1$

■  $Mach = 0.2$

■  $U_{inf} = 0.2$

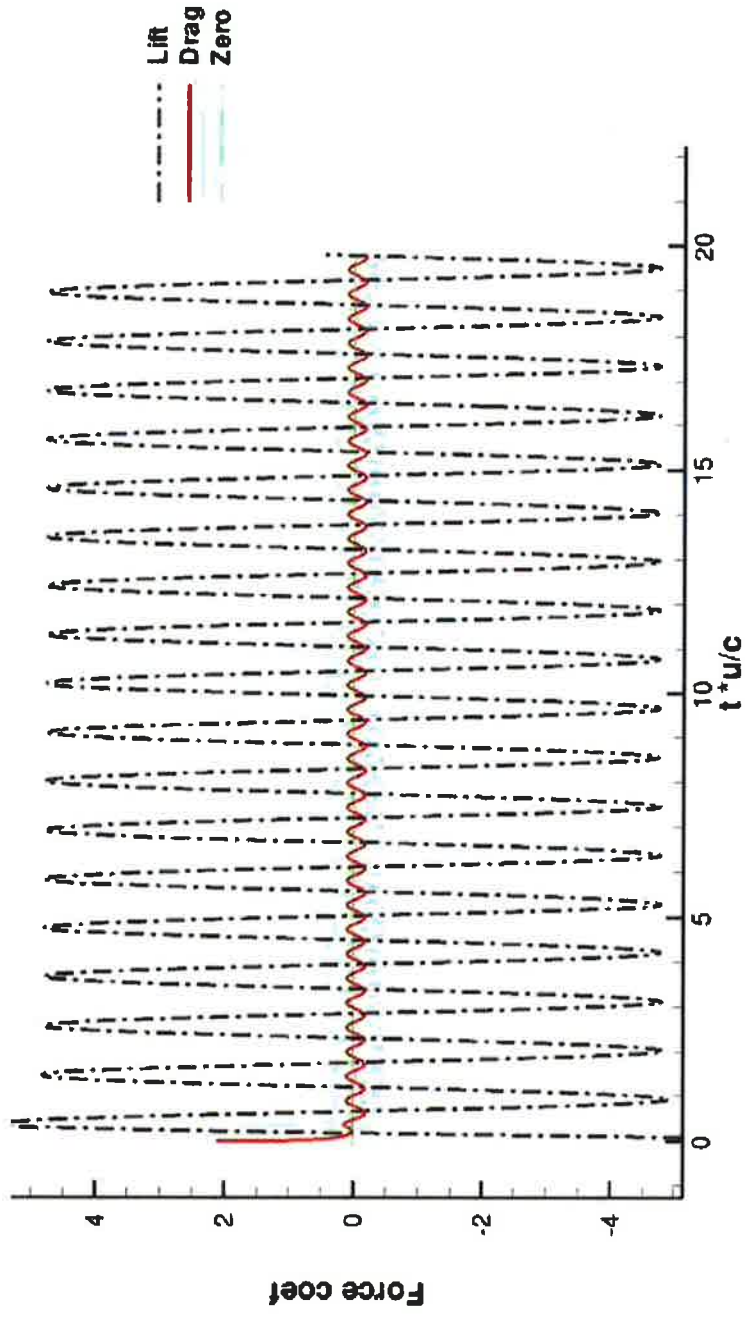
■  $Re = 1850$

■ 4th order spectral difference method



$$-6 \leq \omega c / U_{\infty} \leq 6$$

# Force coefficients for the slowly plunging airfoil



$$-4.78 < C_l < 4.77$$

$$\text{Mean drag coef} = -0.0436$$

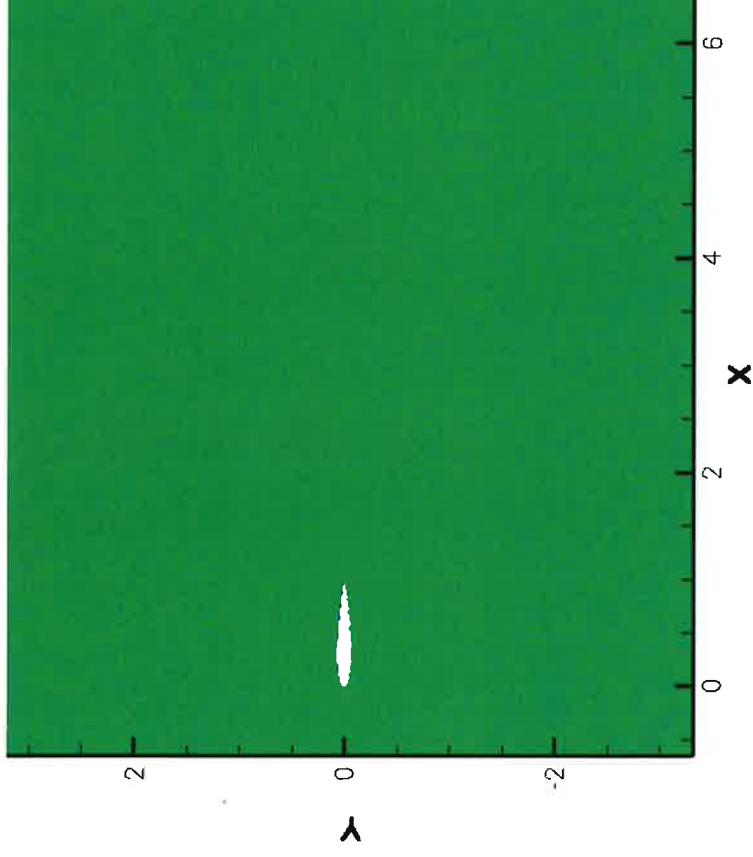
$$-0.181 < C_d < 0.1$$

# Fast plunging case



Our simulations suggest that the asymmetrical flow pattern depends on the direction of the airfoil's first stroke.

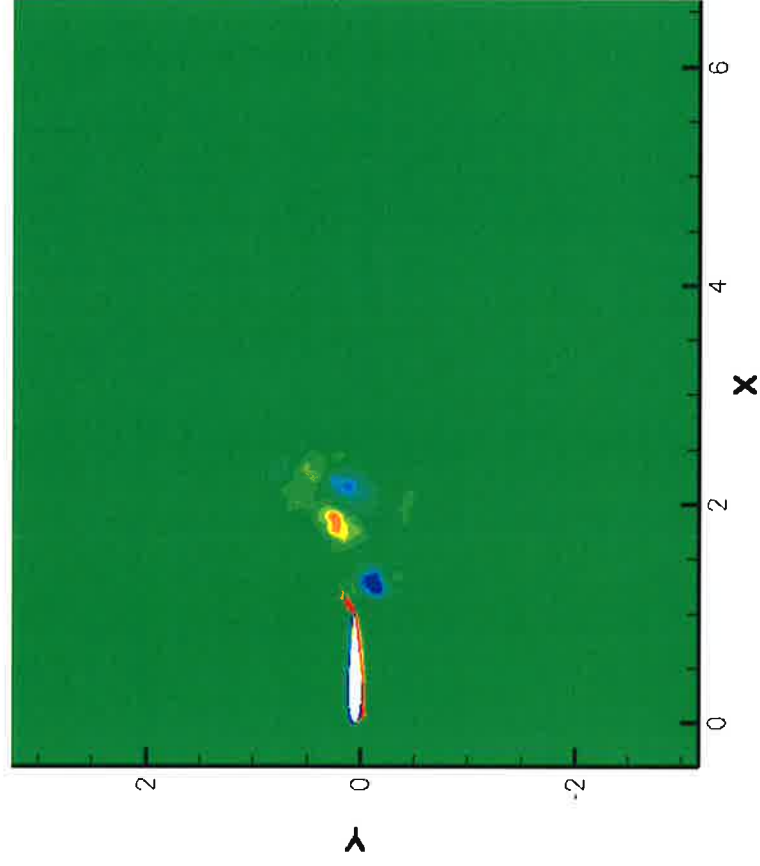
- First stroke of the airfoil goes upwards
- $Y(t) = Ae \cdot \sin(\omega t)$
- $\Omega = 2.46$
- $Ae = 0.12 c$
- Chord  $c = 1$
- $Mach = 0.2$
- $U_{inf} = 0.2$
- $Re = 1850$
- 3<sup>rd</sup> order spectral difference method



# Fast plunging, 3<sup>rd</sup>-order simulation



- First stroke of the airfoil goes downwards
- $Y(t) = Ae \cdot \sin(\omega t)$
- $\Omega = 2.46$
- $Ae = 0.12 c$
- Chord  $c = 1$
- $Mach = 0.2$
- $U_{inf} = 0.2$
- $Re = 1850$
- 3<sup>rd</sup>-order spectral difference method



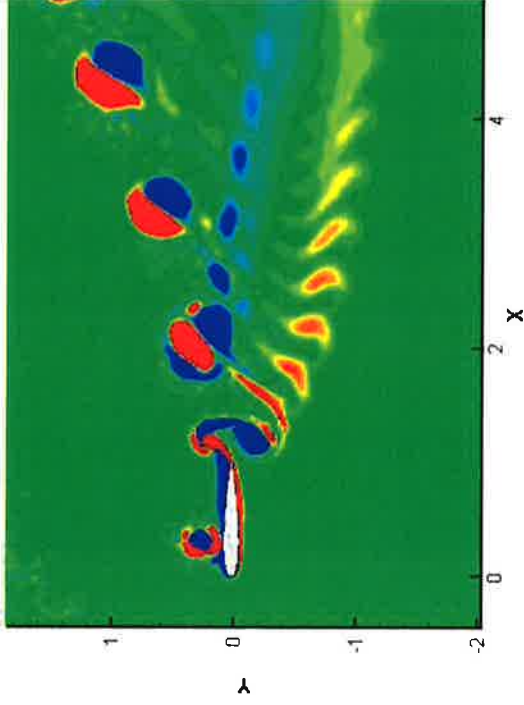




# Simulation versus Experiment



(a) Experiment by Jones, Dohring, and Platzer; AIAA JOURNAL Vol. 36, No. 7, July 1998



(b) 4<sup>th</sup> order Spectral difference method

**Our prediction agrees better with the experimental results than any other published results !**

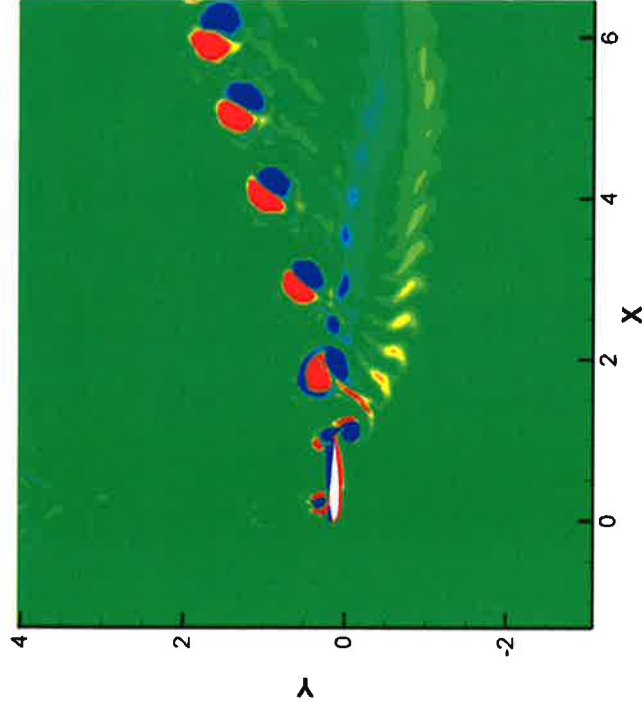
- First stroke of the airfoil goes downwards
- $Y(t) = Ae \cdot \sin(\omega \cdot t)$
- $\Omega = 2.46$
- $Ae = 0.12 c$

- Mach = 0.2
- $U_{inf} = 0.2$
- Re = 1850
- Chord  $c = 1$

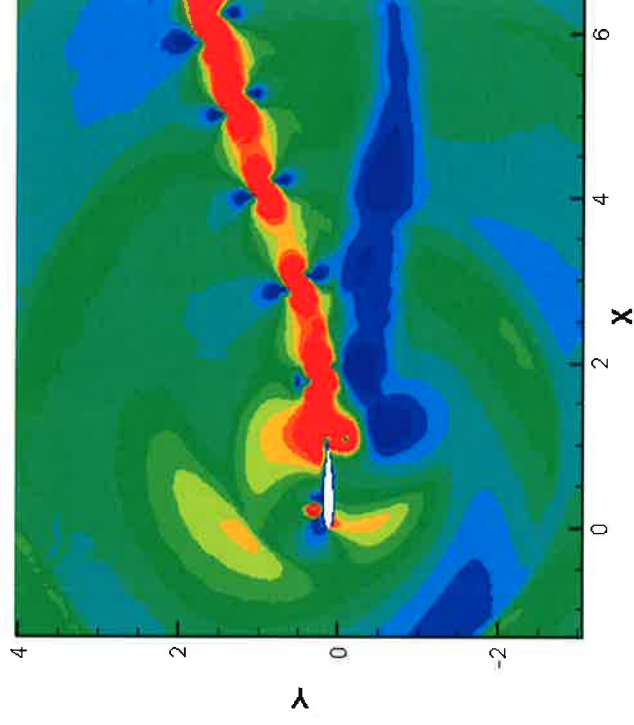




# 4th-order SD prediction for plunging airfoil

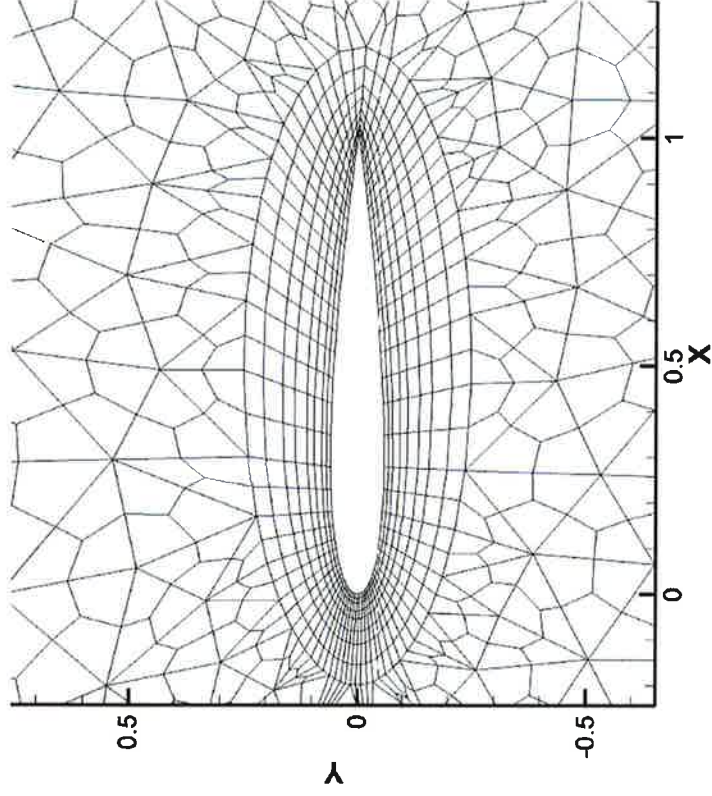


- Vorticity:  $-6 < \Omega * c / U < 6$
- $Y(t) = A_e * \sin(\omega * t)$
- $\Omega = 2.46$
- $A_e = 0.12 c$



- Normalized velocity magnitude:  
 $0.5 < |V| / U_{\infty} < 2$

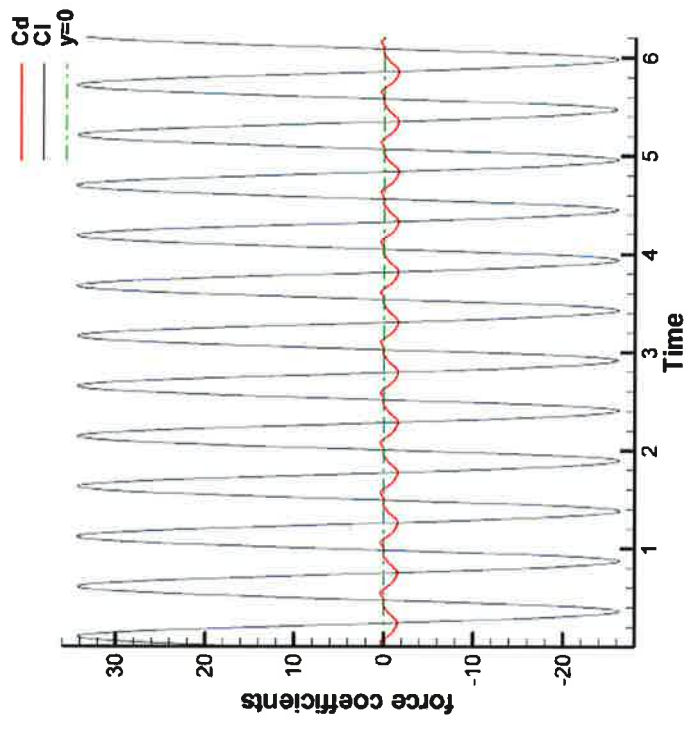
# Force coefficients for fast plunging case



(a) Mesh B

$$-26 < Cl < 34$$

$$-1.51 < Cd < 0.37$$

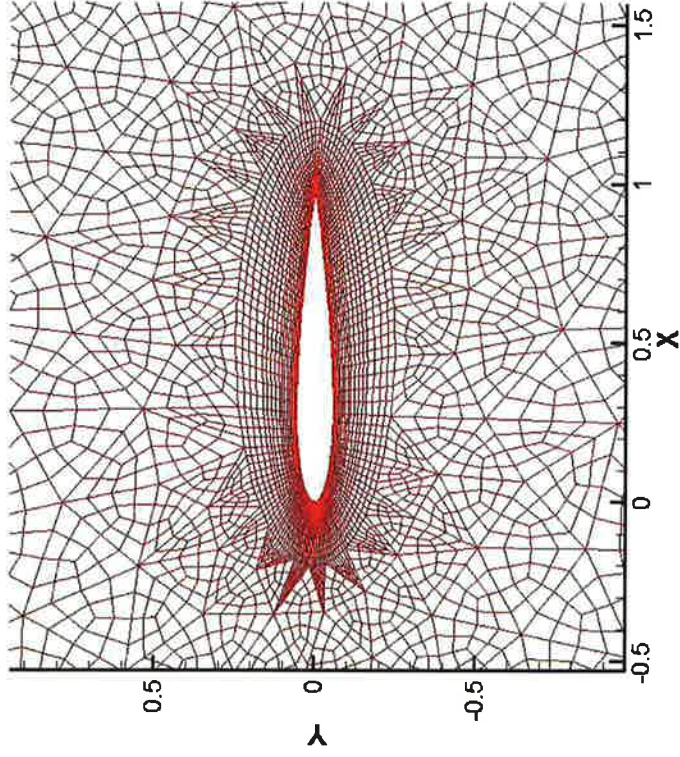


(b) the 4<sup>th</sup> order SD method;

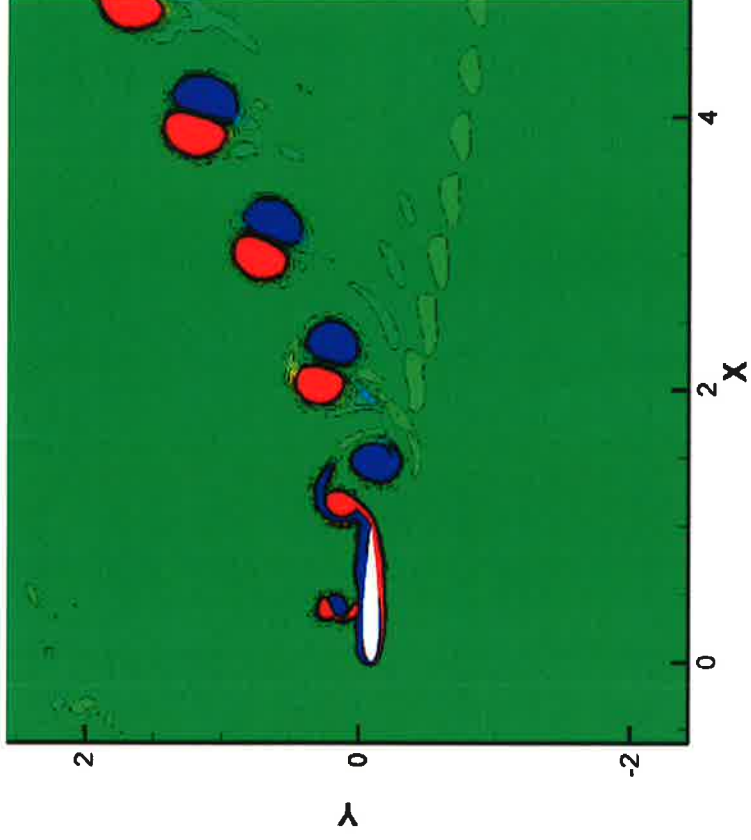
$$\text{Mean drag} = -0.51$$

$$\text{Mean lift} = 2.57$$

# Vorticity shedding on the finest mesh



(a) Mesh C



(b) Vorticity  $-6 \leq \omega c/U_\infty \leq 6$

Third-order SD method is used



## Conclusions



- The SD method is robust and able to model compressible viscous flow with moving boundaries with high accuracy.
- It enables accurate simulations of vortex dominated flows
- Potential applications include flapping wing flight and wind turbine aerodynamics