High-order Spectral Difference Method for Simulating Unsteady Flow past a plunging airfoil

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Aeronautics and Astronautics

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Vortex dominated flow workshop
National Institute of Aerospace
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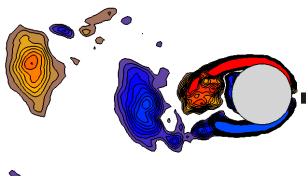
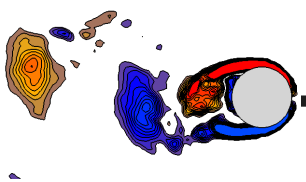


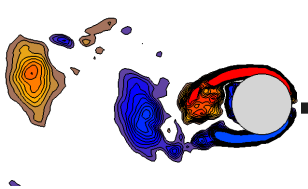
OVERVIEW

- **Introduction**
- **Spectral Difference method**
- **Validation**
- **Vortex Shedding**
- **Subsonic Flow past a plunging airfoil**



Higher order Methods (DG,SV,SD)

- Low numerical dissipation
- Less numerical dispersion
- Unstructured grids
- CPU efficient/easy to parallelize
- Not memory intensive
- Easy to program, universal construction by placing unknown points in a geometrical similar manner
- SD attains a simpler form and higher efficiency than DG and SV.
- Flexible (grid-independence, hp adaptation, moving boundary, deformable grid)



History of Spectral Difference Method

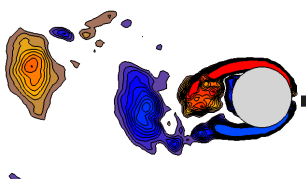
Two early papers on the SD method

- D. A. Kopriva and J. H. Kolias, A conservative staggered-grid chebyshev multidomain method for compressible flows. J. Comput. Phys. 125 (1996), p. 244. **STRUCTURED STAGGERED GRID**
- Yen Liu , Marcel Vinokur , Z. J. Wang, Spectral difference method for unstructured grids I: basic formulation, Journal of Computational Physics, v.216 p.780-801,2006. **UNSTRUCTURED MULTIVARIATE FORMULATION**

Our papers on the SD method

- Wang, Z.J., Liu, Y., May, G., Jameson, A.: "Spectral Difference Method for Unstructured Grids II: Extension to the Euler Equations" , J. Sci. Comput. 32 (1) pp. 54-71, July, 2007. EXTENSITION TO EULER EQUATIONS
- C. Liang, S. Premasuthan, A. Jameson, "High-order accurate simulation of flow past two side-by-side cylinders with Spectral Difference method", 2009, vol 87, pp. 812-817, Journal of Computers and Structures. EXTENSITION TO 2D Viscous flow on quadrilateral elements
- C. Liang, A. Jameson and Z. J. Wang, "Spectral Difference method for two-dimensional compressible flow on unstructured grids with mixed elements", Journal of Computational Physics, vol 228, pp 2847-2858, 2009. EXTENSITION TO 2D Viscous flow on mixed elements

Development of the spectral difference method in Stanford University

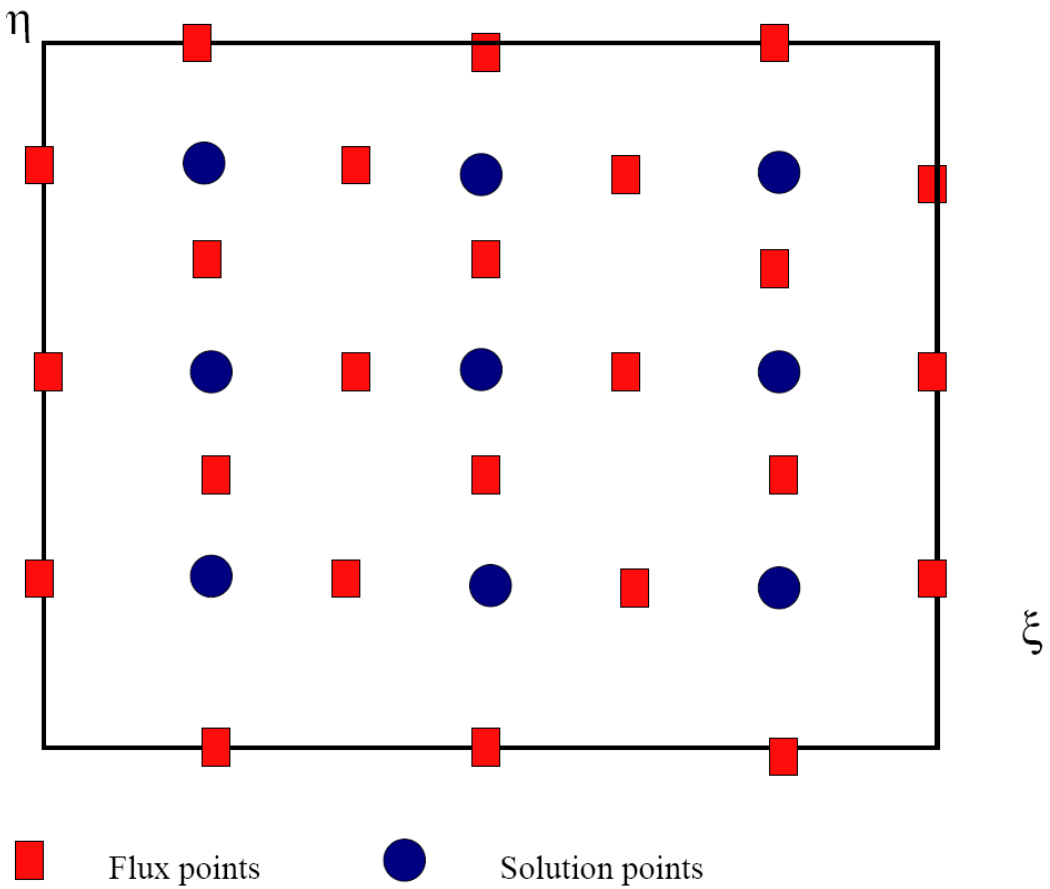


- 2D quadrilateral/triangular element
- 3D hexahedral elements
- 4^{th}-order and higher on unstructured grids
- Implicit LU-SGS time stepping
- P-multigrid method
- Moving and deforming grids
- Artificial viscosity for shock capturing
- Parallelization (MeTis/MPI) of 3D code

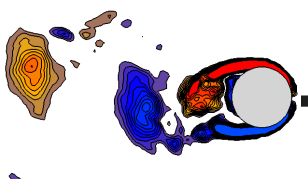
High-order spectral difference method

- dq/dt+df(q)/dx =0

- **The figure shows one way to locally construct a third-order spectral difference scheme for a grid cell.**
- **9 solution points are used.**
- **24 flux points are employed.**
- **The reconstructed field using polynomials is continuous within the cell but discontinuous across the cell boundaries**



- Flux points store f(q)
- Solution points store q



Choosing solution and flux points

- N solution points (one can choose arbitrarily)

(a) $$X_s = \frac{1}{2}\left[1 - cos\left(\frac{2s-1}{2N}\cdot\pi\right)\right], s = 1, 2, \cdots, N.$$ Gauss points

- N+1 flux points

(b) $$X_{s+1/2} = \frac{1}{2}\left[1 - cos\left(\frac{s}{N}\cdot\pi\right)\right], s = 0, 1, \cdots, N.$$ Chebyshev-Gauss-Lobatto points

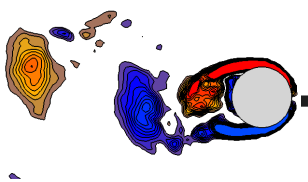
Or

(c) Roots of the equation: $P_n(\xi) = 0$

Legendre-Gauss quadrature points and two end points of 0 and 1.

$$P_{-1}(\xi) = 0$$

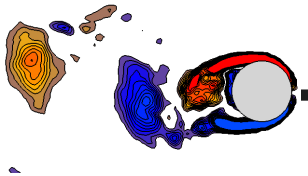
$$P_0(\xi) = 1 \qquad P_n(\xi) = \frac{2n-1}{n}(2\xi - 1)P_{n-1}(\xi) - \frac{n-1}{n}P_{n-2}(\xi)$$



Stability of the Spectral Difference method

- K Abeele, C Lacor and Z. J. Wang (2008), *On the stability and accuracy of the spectral difference method*, Journal of Scientific Computing, vol 37, pp. 162-188.
- Antony Jameson, (2009), *A proof of the stability of the spectral difference method for all orders of accuracy*, Aerospace Computing Lab report, ACL-2009-1, Stanford University, March 2009.
- Both confirm stability when the interior flux points are at the zeros of the Legendre polynomial

Solution and Flux reconstruction



Degree N-1 polynomial through solution points using Lagrange basis polynomial

$$h_i(X) = \prod_{s=1, s\neq i}^{N} \left(\frac{X - X_s}{X_i - X_s} \right)$$

Degree N polynomial through flux points using Lagrange basis polynomial

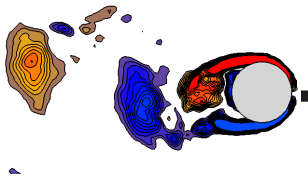
$$l_{i+1/2}(X) = \prod_{s=0, s\neq i}^{N} \left(\frac{X - X_{s+1/2}}{X_{i+1/2} - X_{s+1/2}} \right)$$

Reconstructed solution as well as flux polynomials within an element are written as tensor products of three 1D polynomials

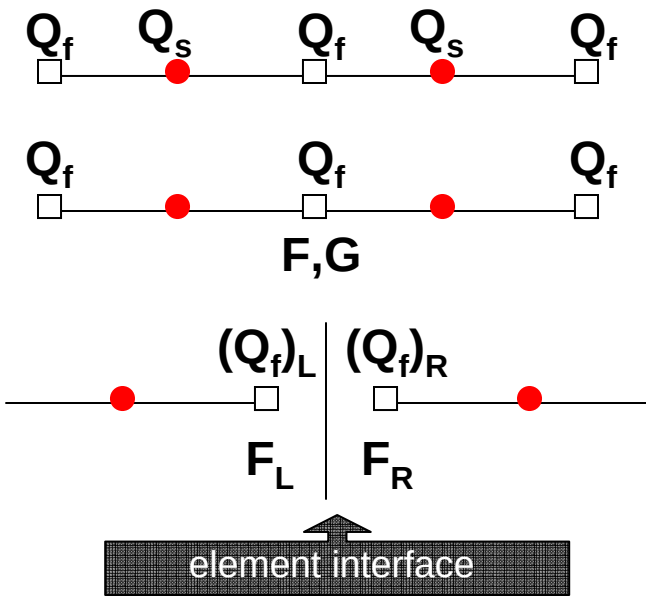
$$Q(\xi, \eta, \beta) = \sum_{k=1}^{N} \sum_{j=1}^{N} \sum_{i=1}^{N} \frac{\tilde{Q}_{i,j,k}}{|J_{i,j,k}|} h_i(\xi) \cdot h_j(\eta) \cdot h_k(\beta)$$

$$\tilde{F}(\xi, \eta, \beta) = \sum_{k=1}^{N} \sum_{j=1}^{N} \sum_{i=0}^{N} \tilde{F}_{i+1/2,j,k} l_{i+1/2}(\xi) \cdot h_j(\eta) \cdot h_k(\beta)$$

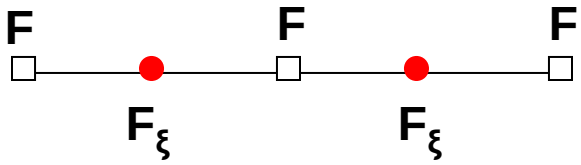
Calculation of inviscid flux derivatives



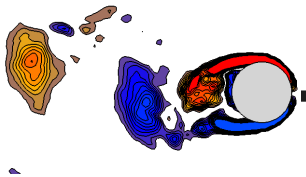
- Given the conservative variables at the solution points, the conservative variables are extrapolated to the flux points
- The inviscid fluxes at the interior flux points are computed
- The inviscid fluxes at the element interfaces are computed using an approximate Reimann solver
- The derivative of the fluxes are computed at the solution points according to



$$\left(\frac{\partial \tilde{F}}{\partial \xi}\right)_{i,j} = \sum_{r=0}^{N} \tilde{F}_{r+1/2,j} \cdot l'_{r+1/2}(\xi_i)$$



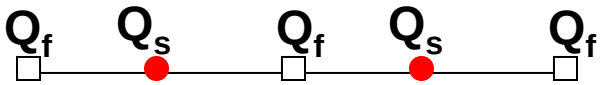
Calculation of viscous flux derivatives (1)



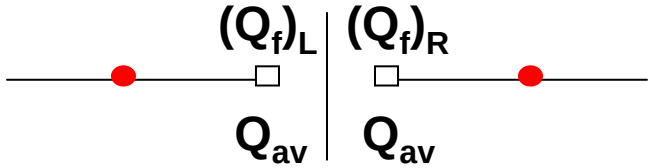
- The viscous fluxes depend on both the solution and the solution gradient at the flux points

$$F_v = F_v(Q_f, \nabla Q_f)$$

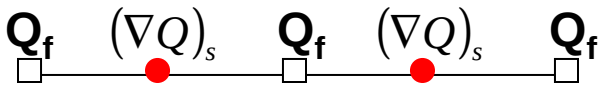
- Reconstruct conservative variables at flux points using values at solution points



- At flux points on element interfaces, compute the average of left and right solution values

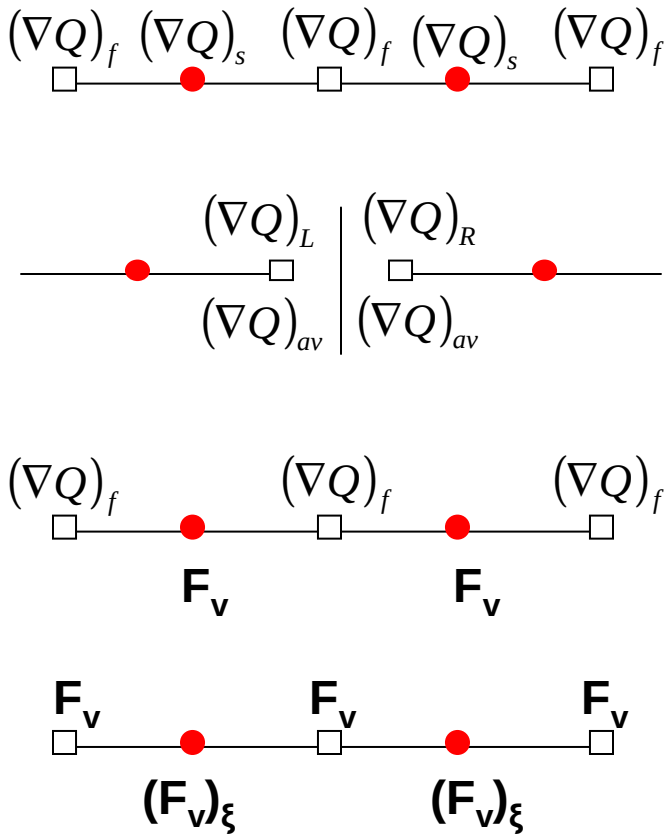


- Calculate the gradient of Q at the solution points from Q at flux points



Calculation of viscous flux derivatives (2)

- Reconstruct ∇Q from solution points to flux points
- At flux points on element interfaces, compute the average of left and right ∇Q values
- Compute viscous fluxes at flux points
- Compute viscous flux derivatives at solution points
- We are ready to compute the NS equation, using a five-stage R-K scheme

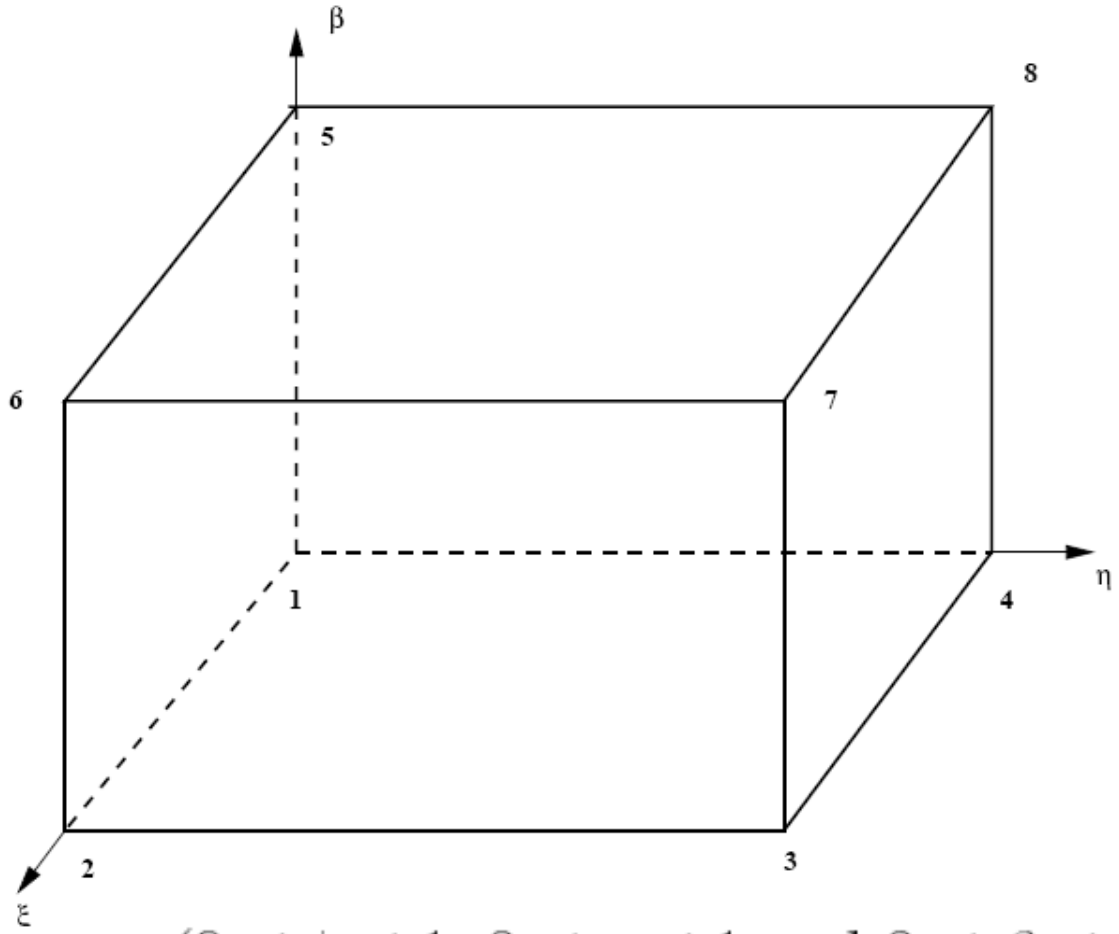


$$\frac{\partial Q}{\partial t} + \nabla F_e(Q) - \nabla F_v(Q, \nabla Q) = 0$$

Hexahedral grid mapping

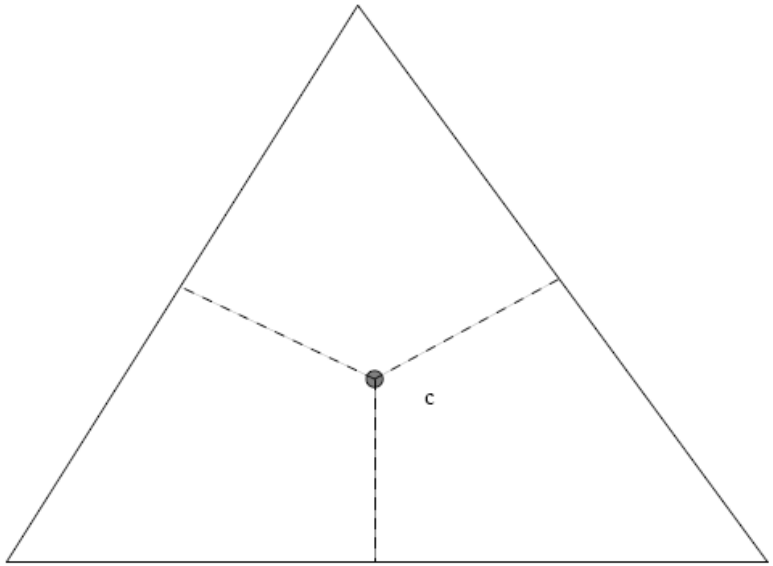
$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \sum_{i=1}^{K} M_i(\xi, \eta, \beta) \begin{pmatrix} x_i \\ y_i \\ z_i \end{pmatrix}$$

- Mapping from physical hexahedral grid elements to a standard computational element

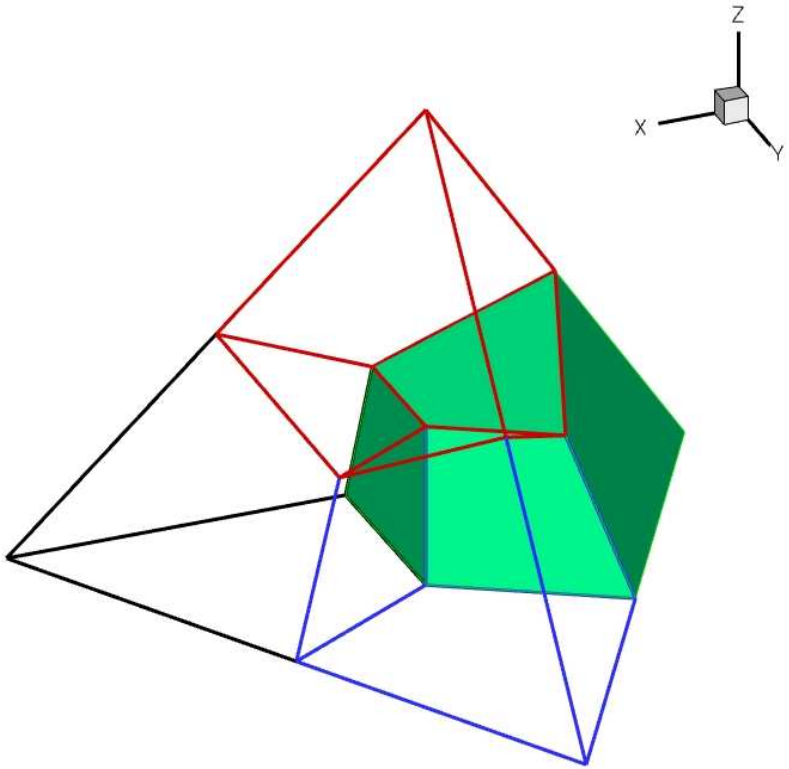


$$(0 \le \xi \le 1,\ 0 \le \eta \le 1 \text{ and } 0 \le \beta \le 1)$$

Splitting triangulated grids to quadrilateral grids in 2D and 3D



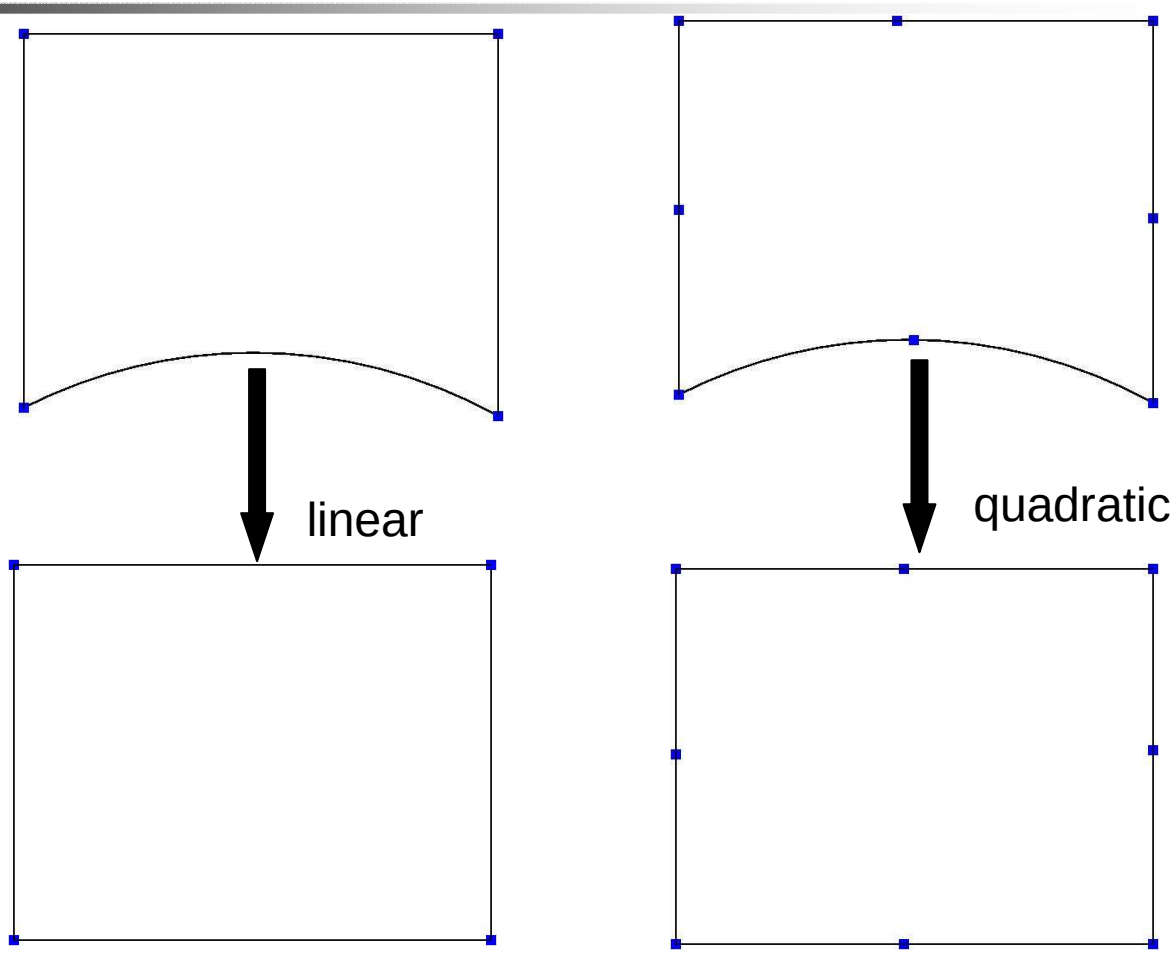
(a) A triangle split into three quadrilaterals



(b) A tetrahedron split into four hexahedra

Curved wall representation

$$\begin{pmatrix} x \\ y \end{pmatrix} = \sum_{i=1}^{K} M_i(\xi,\eta) \begin{pmatrix} x_i \\ y_i \end{pmatrix}$$



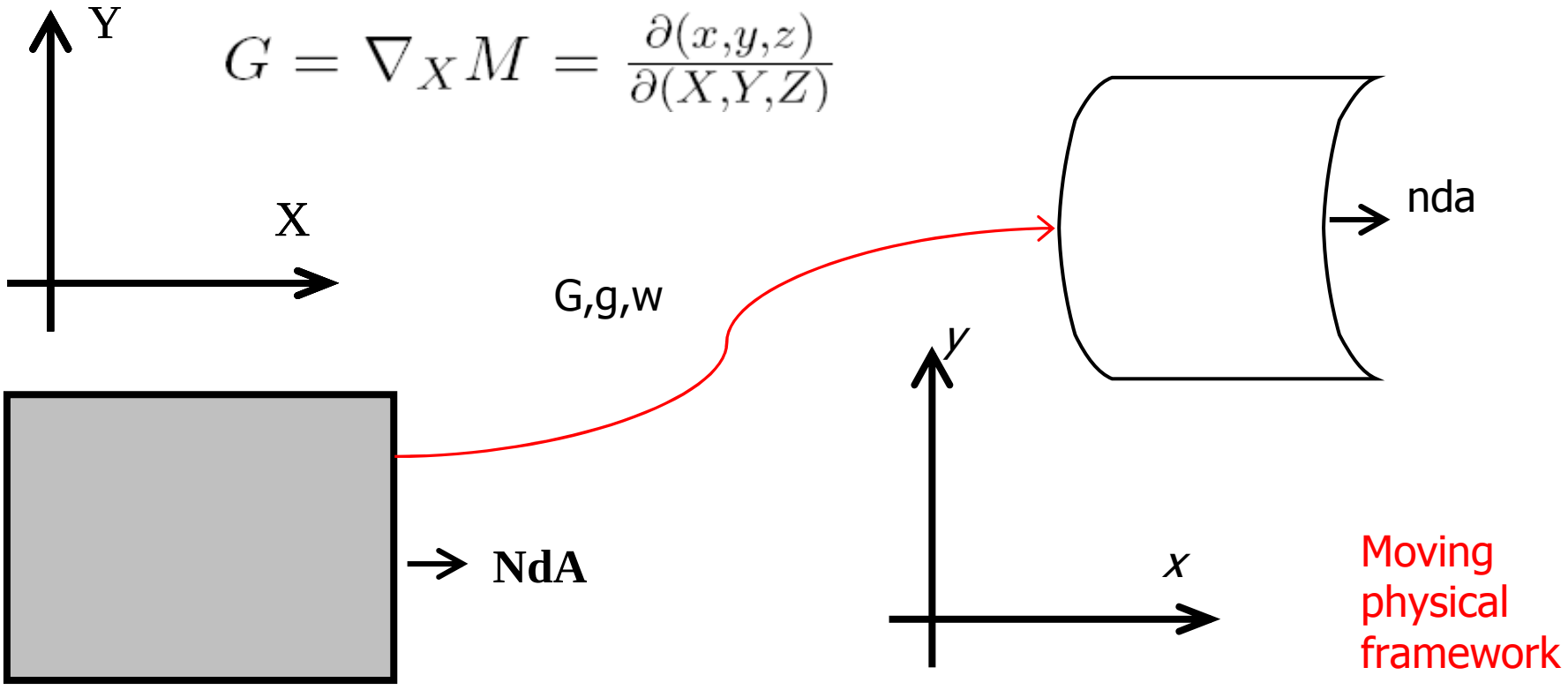
M is shape function; K =8, 20 for 3D linear, quadratic mapping

K=4,8 for 2D linear, quadratic mapping

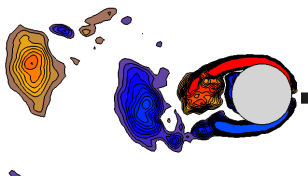
Extension to Moving and Deformable Grids (I)

$$(x, y, z) = M(X, Y, Z, t)$$

$$G = \nabla_X M = \frac{\partial(x,y,z)}{\partial(X,Y,Z)}$$



Extension to Moving and Deformable Grids (II)



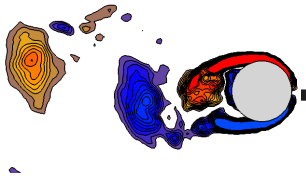
Define a mapping velocity :

$$\frac{\partial M(X,Y,Z,t)}{\partial t} = w$$

$$\frac{\partial \hat{Q}}{\partial t}|_X + \nabla_X \hat{F}_e(\hat{Q}) - \nabla \hat{F}_v(\hat{Q}, \nabla \hat{Q}) = 0$$

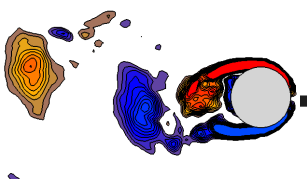
$$g = \det(\boldsymbol{G}) \qquad \hat{Q} = gQ$$

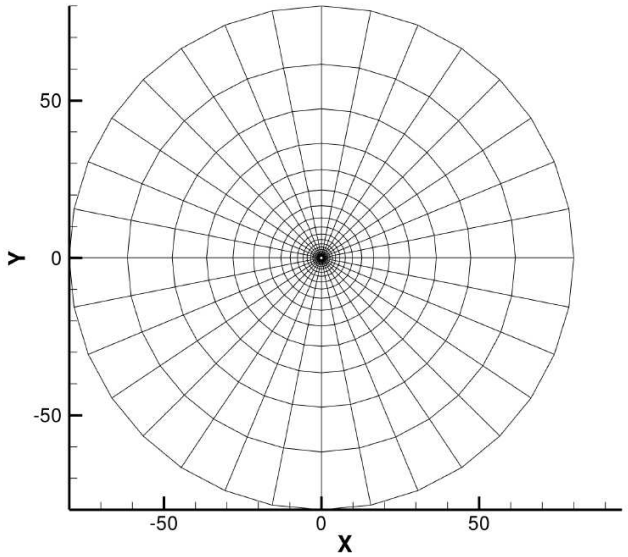
$$\hat{F} = gG^{-1}F - \hat{Q}G^{-1}w$$



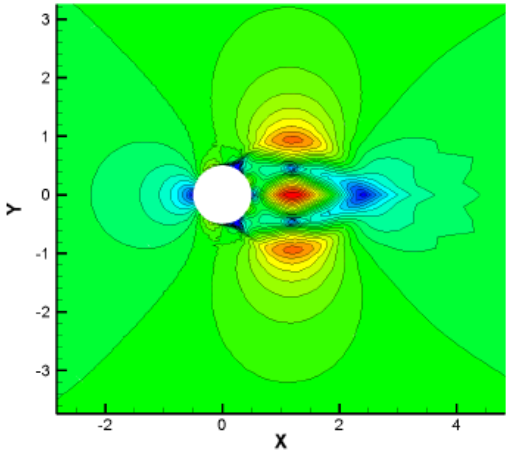
VALIDATION

Mach contour around a circle

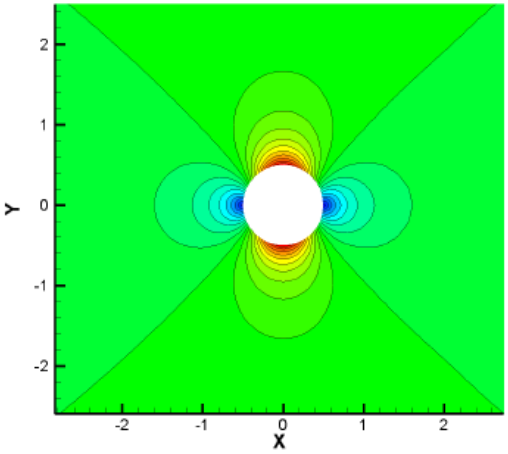




4th-order SD method on a grid with 32x32 cells

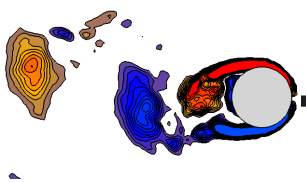


(a) linear wall boundary



(b) cubic wall boundary

- Free stream Mach = 0.2
- Inviscid Euler solver

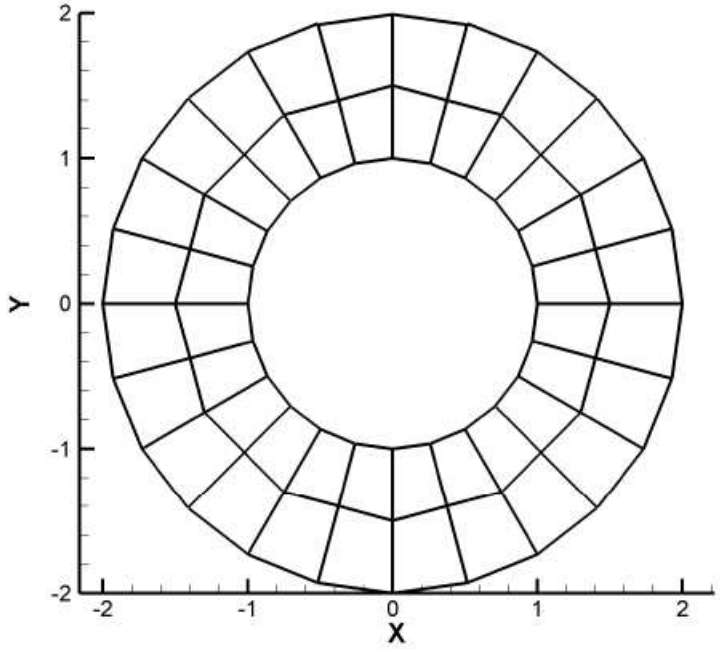


Effect of interface flux on Cd

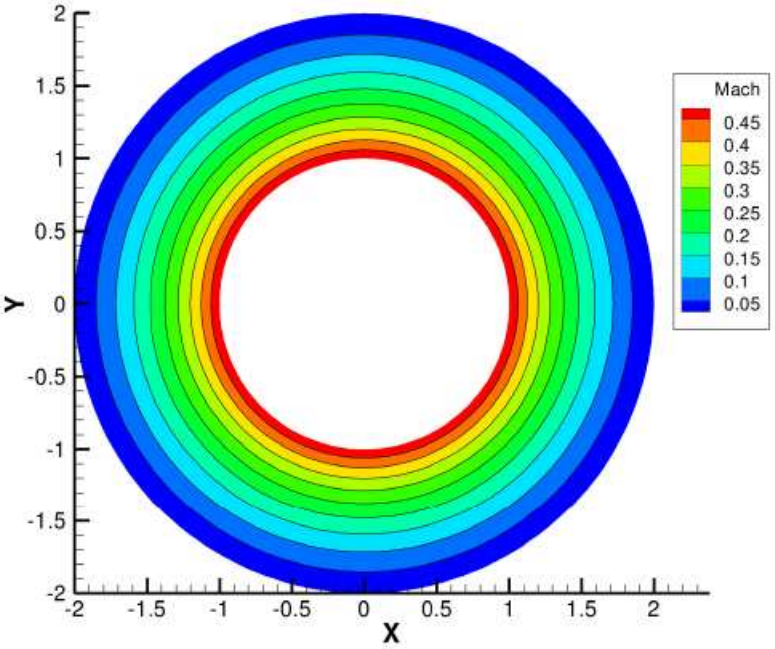
Riemann	order	cell no.	Wall	Mach	$\Delta_t u_\infty / D$	C_d
CUSP	4th	640	quadratic	0.2	2e-4	-1.86e-5
AUSM	4th	640	quadratic	0.2	2e-4	-4.39e-5
Roe	4th	640	quadratic	0.2	2e-4	-1.03e-5
flux vector split	4th	640	quadratic	0.2	2e-4	-1.18e-5
Scalar diffusion	4th	640	quadratic	0.2	2e-4	-8.8e-6

- Inviscid flow past a circle

Compressible Taylor-Couette Flow

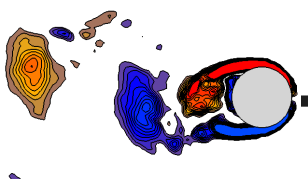


(a) Grid



(b) Mach contour

Mach =0.5, Re=10, isothermal for inner cylinder and adiabatic wall for outer cylinder

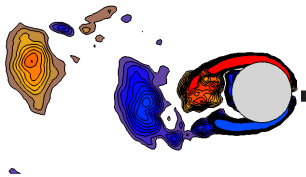


Verification of order of accuracy

No. of elements	No. of DOFs	L2-error	Order
3rd order SD			
48	432	8.896E-04	-
192	1728	1.002E-04	3.15
768	6912	1.084E-05	3.21
4th order SD			
48	768	1.4815E-04	-
192	3072	1.0036E-05	3.88
768	12288	6.5746E-07	3.93

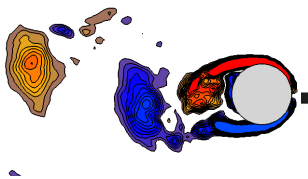
Exact solution for angular velocity

$$v_\theta = \Omega_i r_i \frac{\frac{r_o}{r} - \frac{r}{r_o}}{\frac{r_o}{r_i} - \frac{r_i}{r_o}}$$

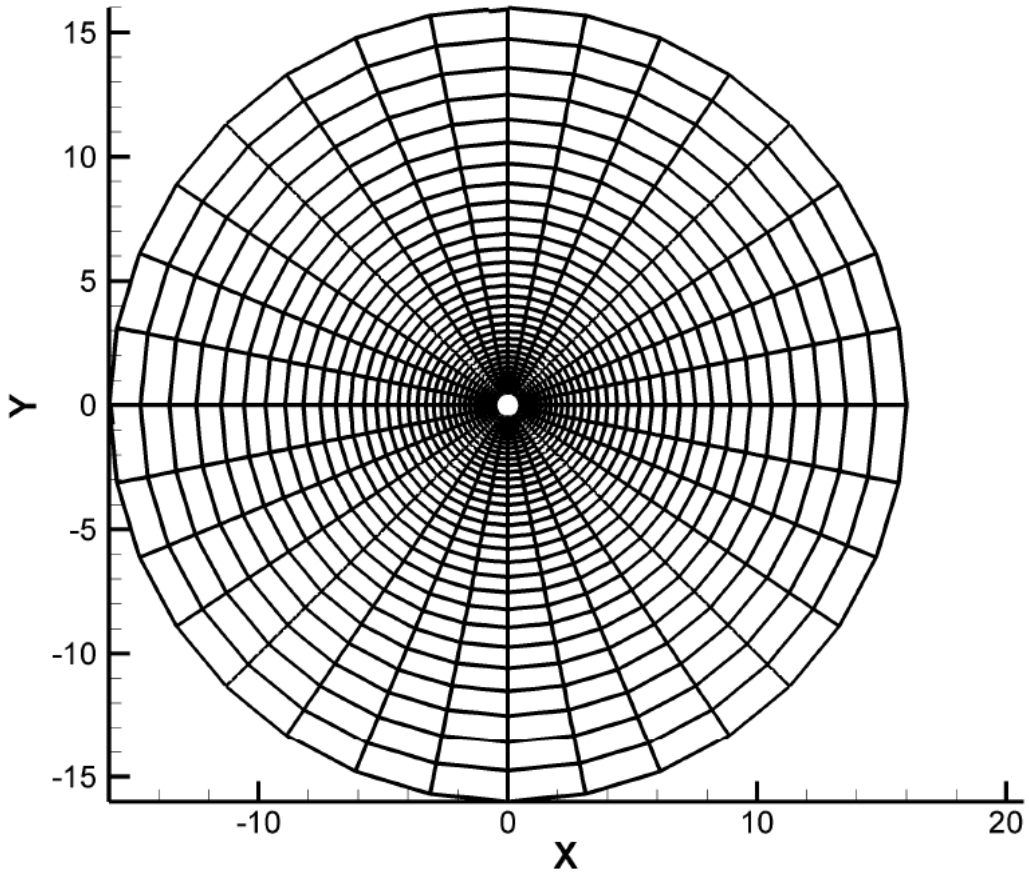


Studies of Vortex Shedding Past a Plunging Cylinder

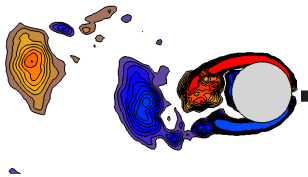
- E. GUILMINEAU and P. QUEUTEY,
A NUMERICAL SIMULATION OF VORTEX SHEDDING FROM AN OSCILLATING CIRCULAR CYLINDER
Journal of Fluids and Structures
Vol 16, 2002, pp. 773-794



Computational grid for oscillating cylinder

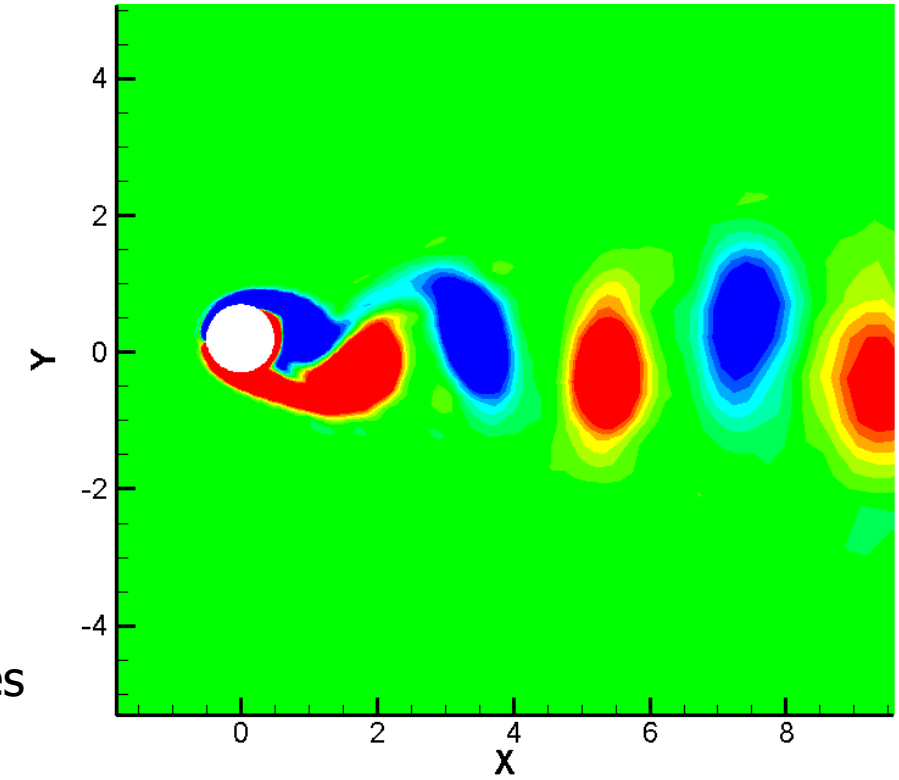


- Our simulation uses only 32x32 grid
- The third-order SD method
- Total degrees-of-freedom = 96x96
- Quadratic curved wall



vorticity shedding behind an oscillating cylinder

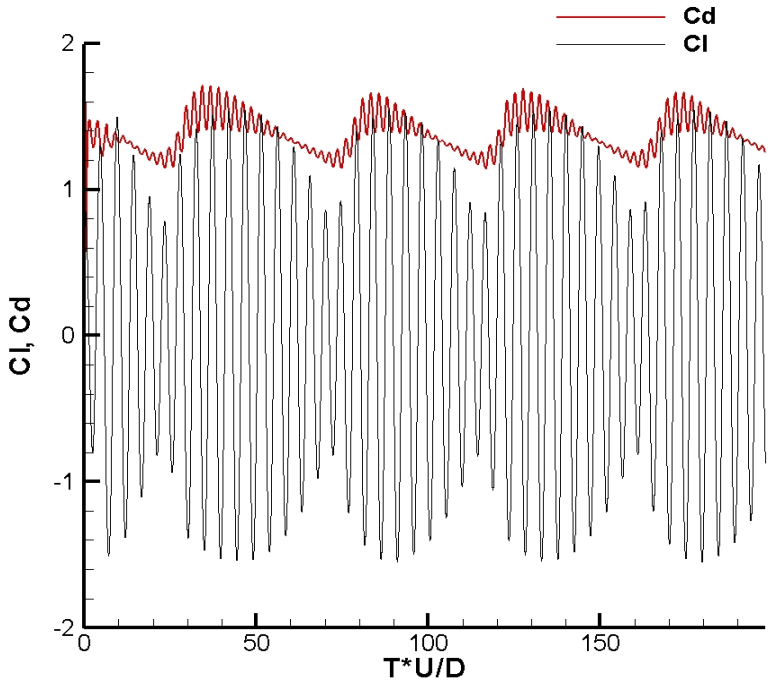
- Y(t)=Ae*cos(2*pi*fe*t) for cylinder transverse motion
- Ae = 0.2D
- fe = 1.1 fn
- -1<omega*d/U<1
- Re =185



- The cylinder motion is identical to the one reported in E. GUILMINEAU and P. QUEUTEY, Journal of Fluids and Structures Vol 16, 2002, pp. 773-794

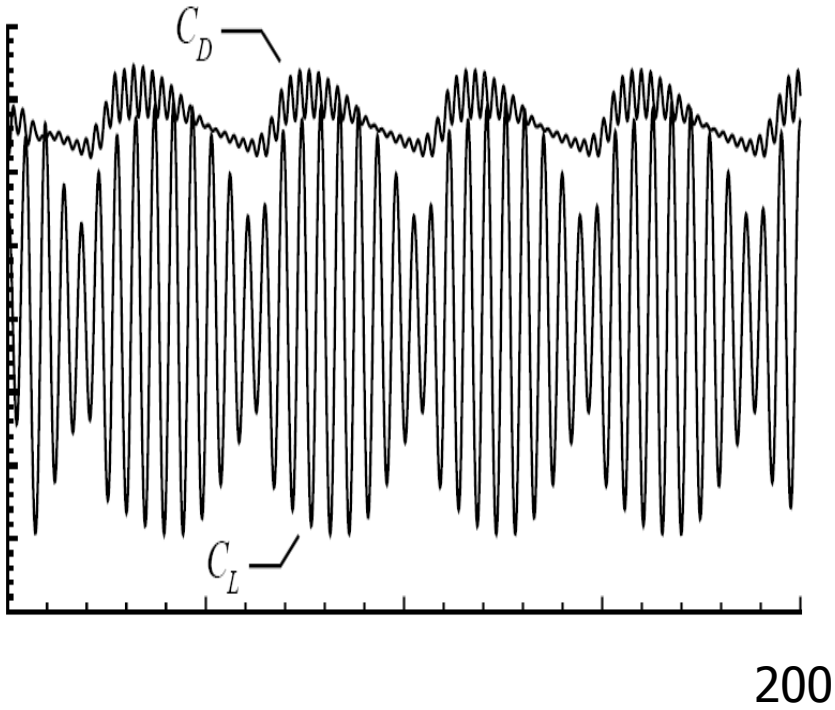
- We obtained nearly identical results using a grid with only 1024 cells compared to their 48,000 cells

Fluid-exerted forces



(a) Third order spectral difference

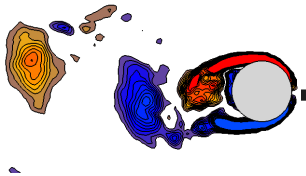
Total cell : 32x32



(b) E. GUILMINEAU and P. QUEUTEY
Journal of Fluids and Structures
Vol 16, 2002, pp. 773-794

Total cell : 240x200

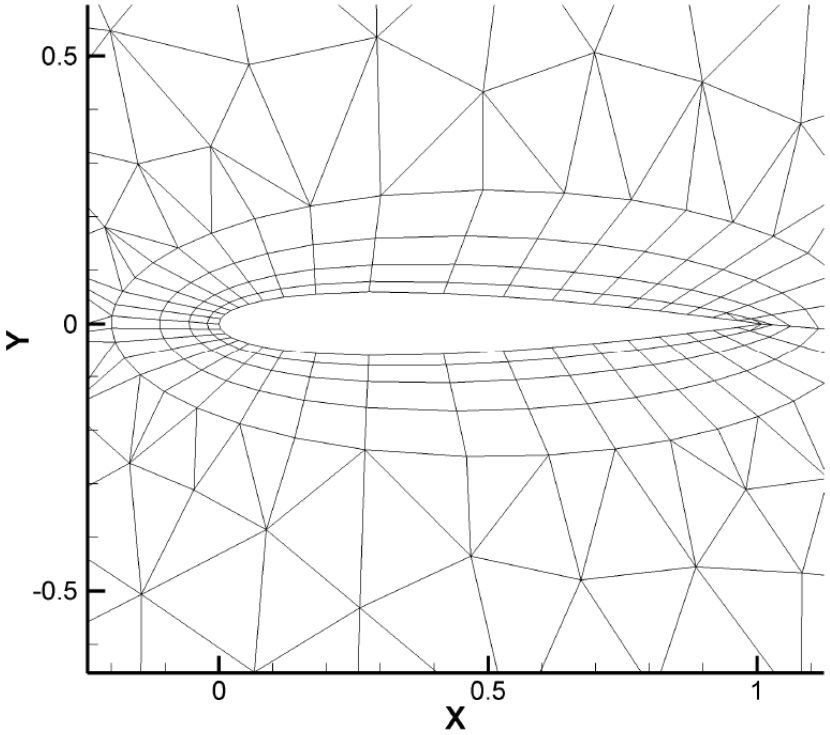
The predicted force coefficients are in a good agreement with Guilmineau and Queutey (2002) and much less computational cells are used.



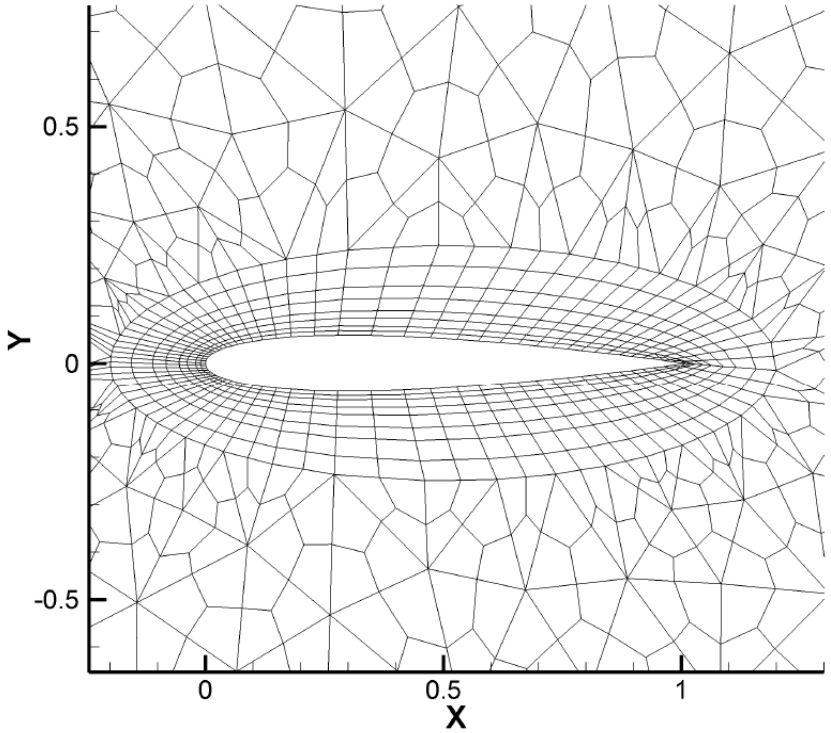
Studies of plunging airfoils

- K. D. Jones, C. M. Dohring and M. F. Platzer, Experimental and computational investigation of the Knoller-Betz effect, *AIAA Journal*, Vol 36, pp. 780-783.

NACA 0012 grid with mixed elements

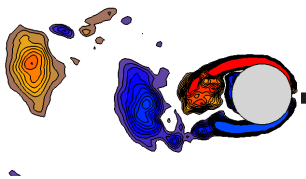


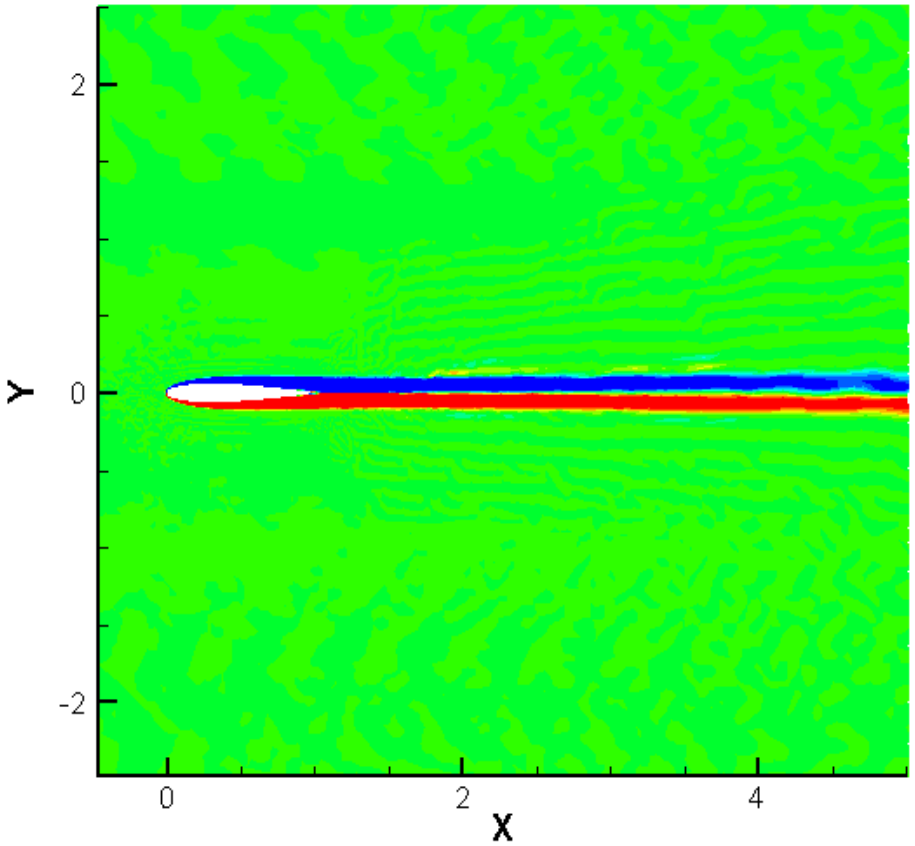
(a) Initial hybrid mesh: 'Mesh A'



(b) One-level h-refinement: `Mesh B'

Viscous flow past a stationary airfoil





- Re = 1850
- Mach = 0.2
- Vortices:

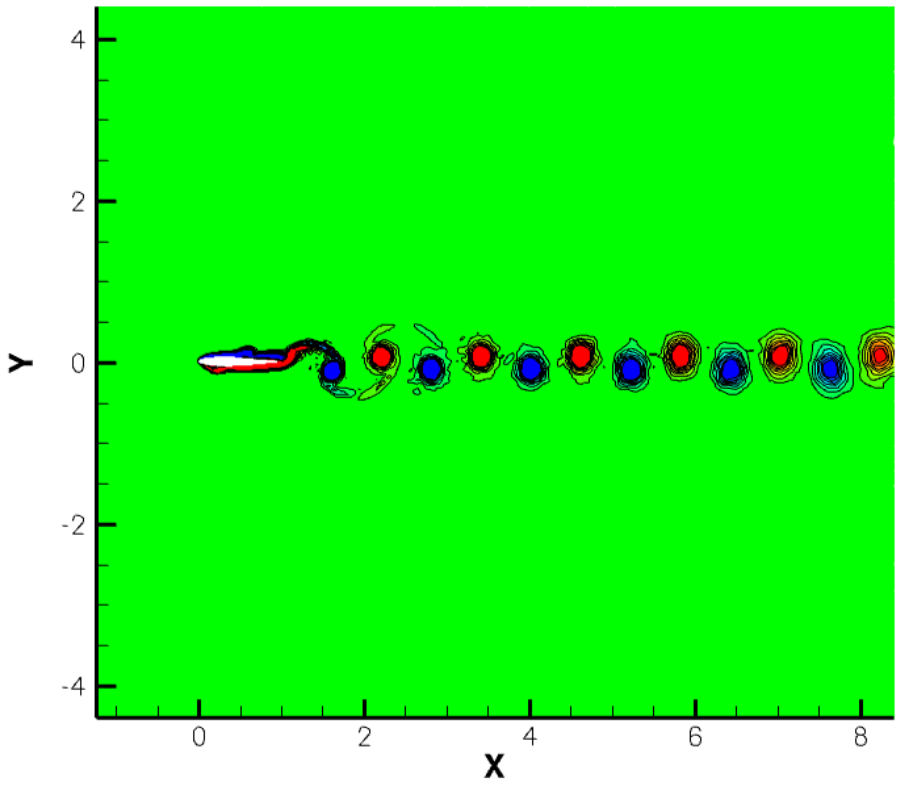
$$-6 \leq \omega c/U_\infty \leq 6$$

Steady flow solution

is reached eventually

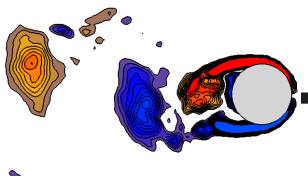
Slow plunging airfoil case

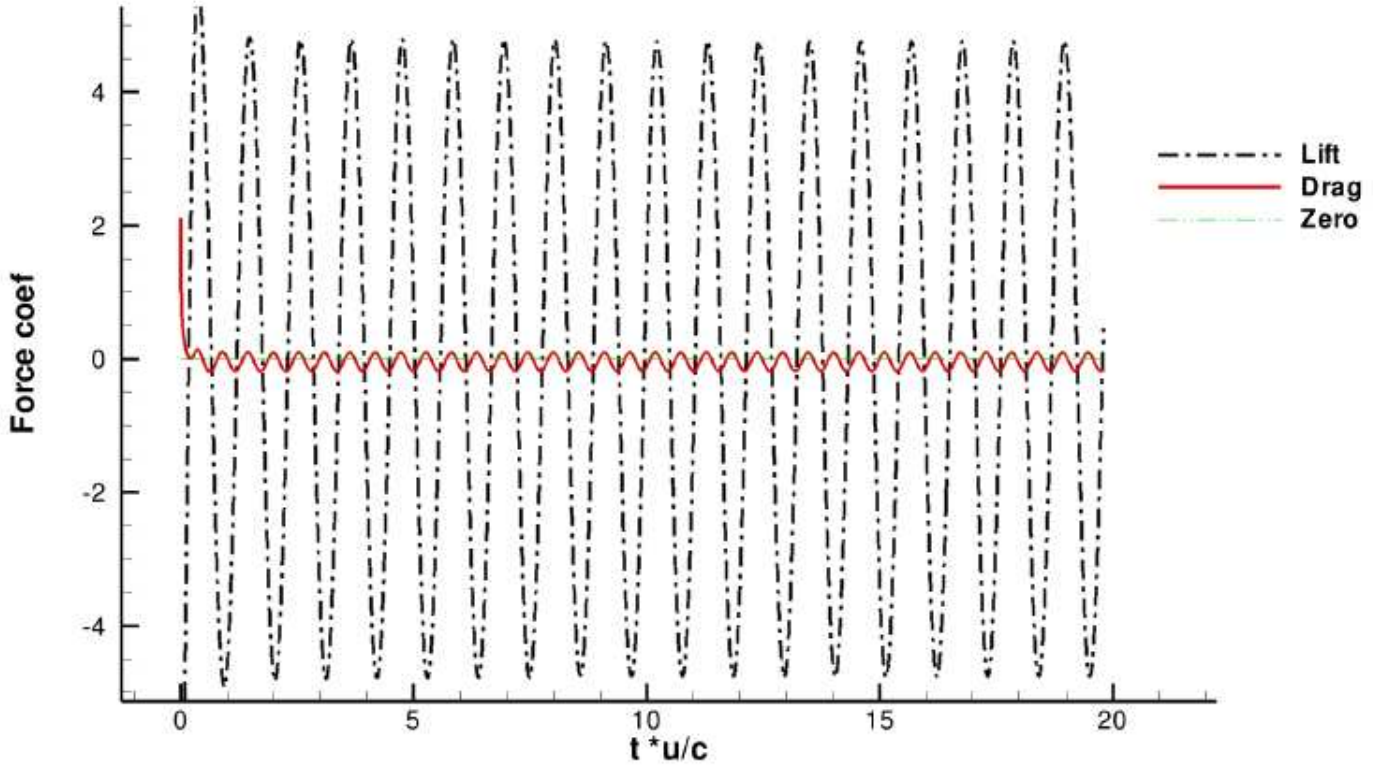
- Y(t)=Ae*sin(omega*t)
- Omega = 1.15
- Ae = 0.08 c
- Chord c =1
- Mach = 0.2
- Uinf = 0.2
- Re=1850
- 4th order spectral difference method



$$-6 \leq \omega c/U_\infty \leq 6$$

Force coefficients for the slowly plunging airfoil





-4.78<Cl<4.77

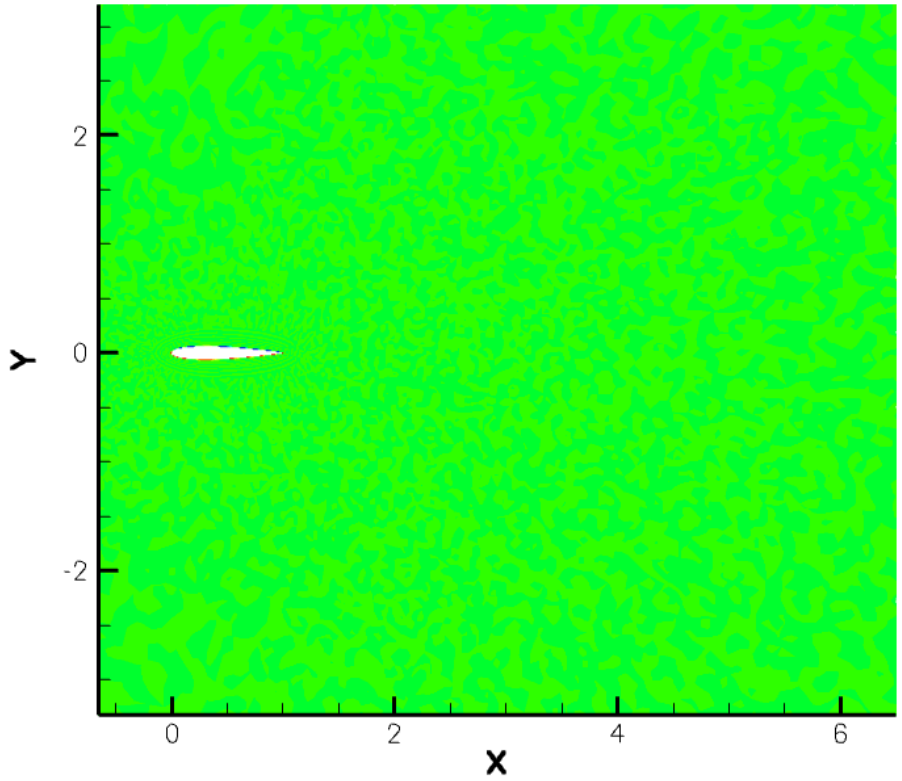
-0.181<Cd<0.1

Mean drag coef = -0.0436

Fast plunging case

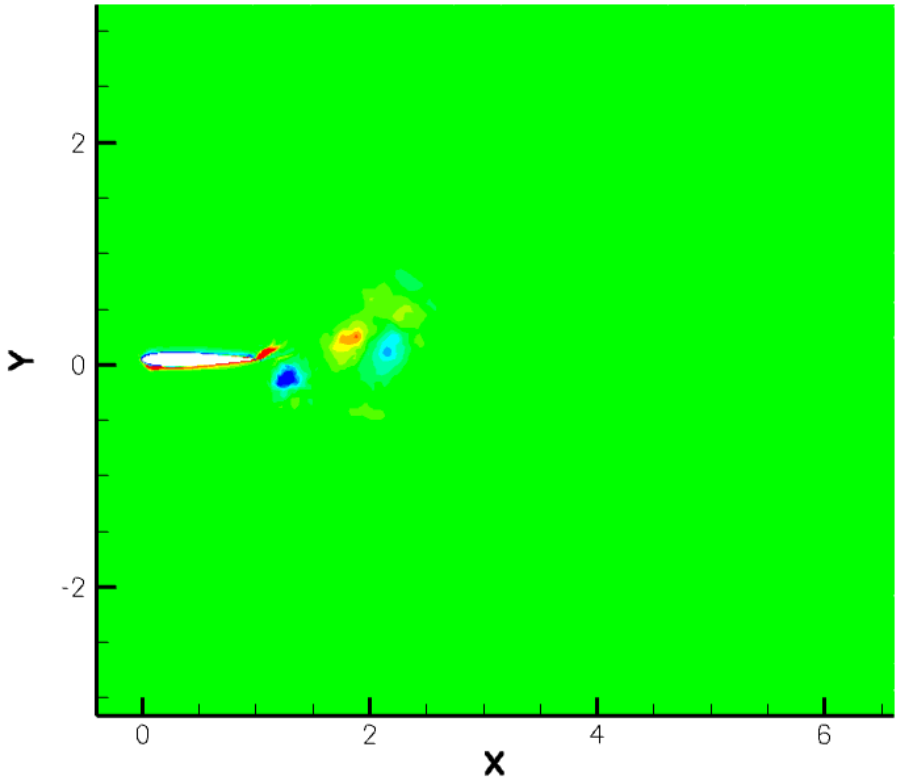
Our simulations suggest that the asymmetrical flow pattern depends on the direction of the airfoil's first stroke.

- First stroke of the airfoil goes upwards
- Y(t)=Ae*sin(omega*t)
- Omega = 2.46
- Ae = 0.12 c
- Chord c =1
- Mach = 0.2
- Uinf = 0.2
- Re=1850
- 3rd order spectral difference method

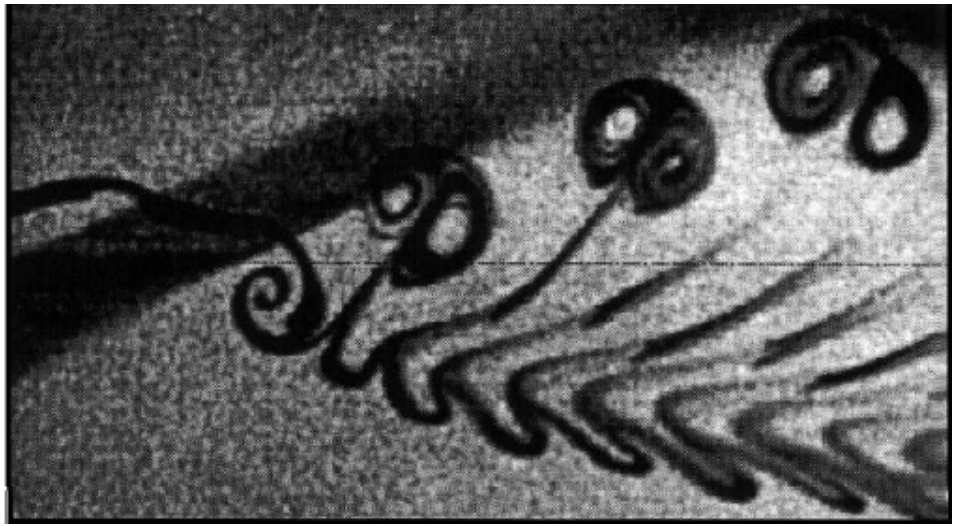


Fast plunging, 3rd-order simulation

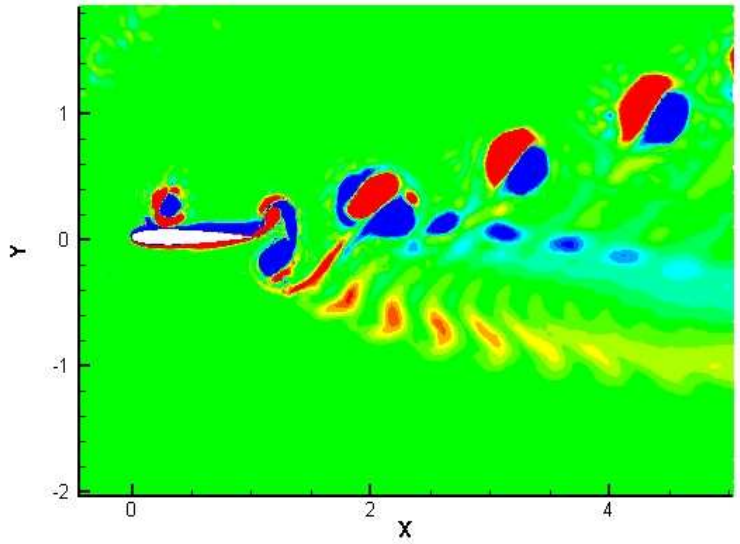
- First stroke of the airfoil goes downwards
- Y(t)=Ae*sin(omega*t)
- Omega = 2.46
- Ae = 0.12 c
- Chord c =1
- Mach = 0.2
- U_inf =0.2
- Re=1850
- 3rd –order spectral difference method



Simulation versus Experiment



(a) Experiment by Jones, Dohring, and Platzer; AIAA JOURNAL Vol. 36, No. 7, July 1998



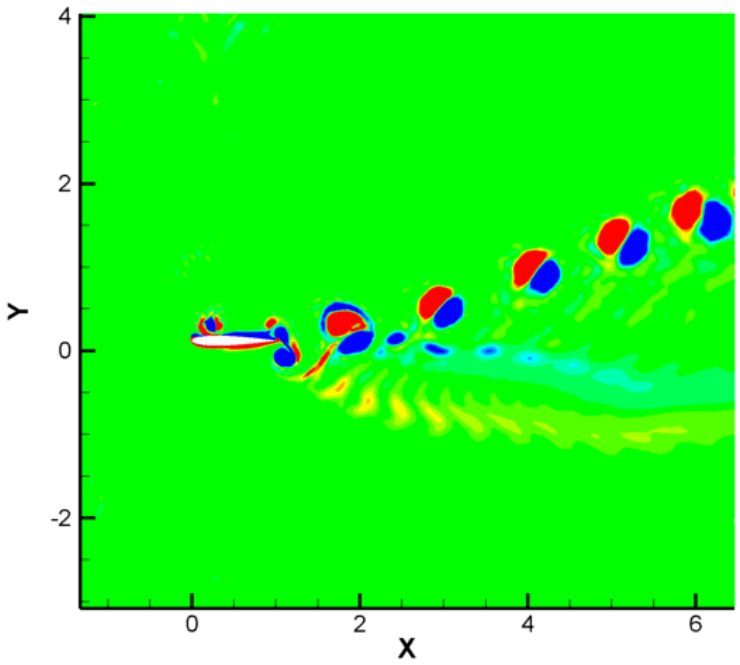
(b) 4th order Spectral difference method

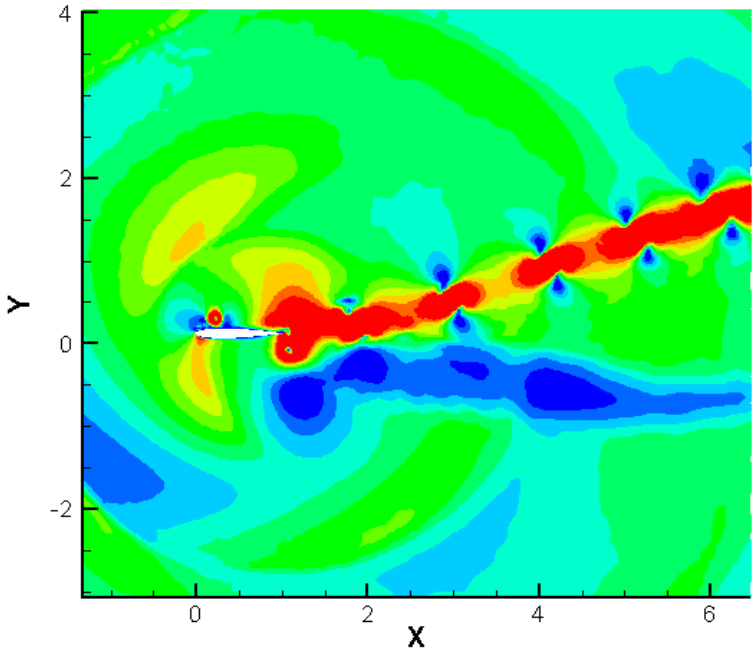
Our prediction agrees better with the experimental results than any other published results !

- First stroke of the airfoil goes downwards
- Y(t)=Ae*sin(omega*t)
- Omega = 2.46
- Ae = 0.12 c

- Mach = 0.2
- Uinf = 0.2
- Re=1850
- Chord c=1

4th-order SD prediction for plunging airfoil



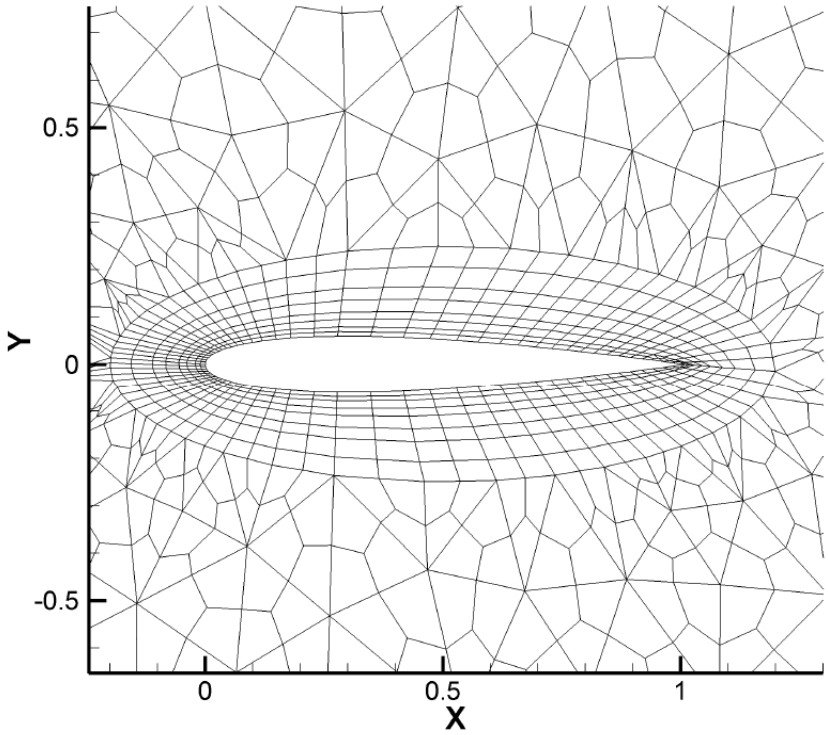


- Vorticity: -6<Omega*c/U<6
- Y(t)=Ae*sin(omega*t)
- Omega = 2.46
- Ae = 0.12 c

- Normalized velocity magnitude:

  0.5<|V|/U_inf<2

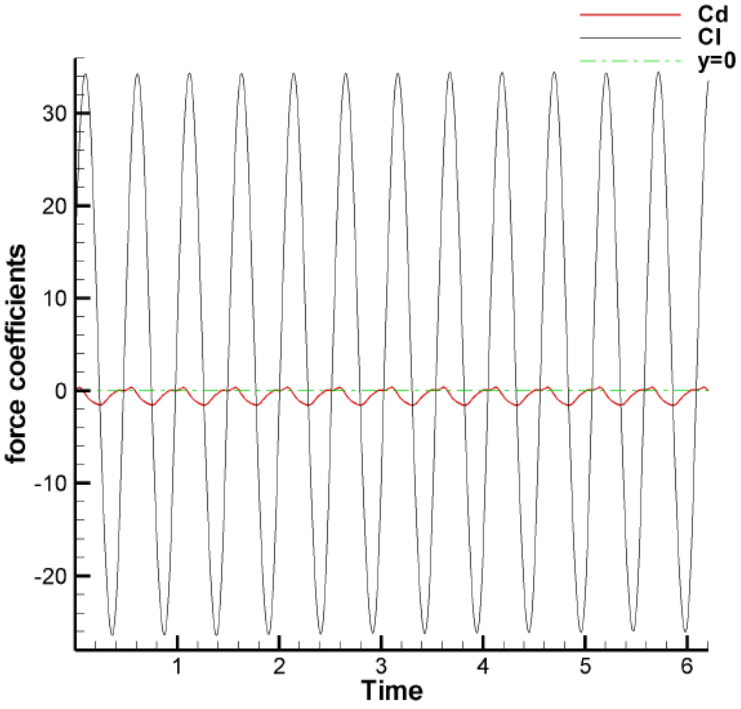
Force coefficients for fast plunging case



(a) Mesh B

-26<Cl<34

-1.51<Cd<0.37

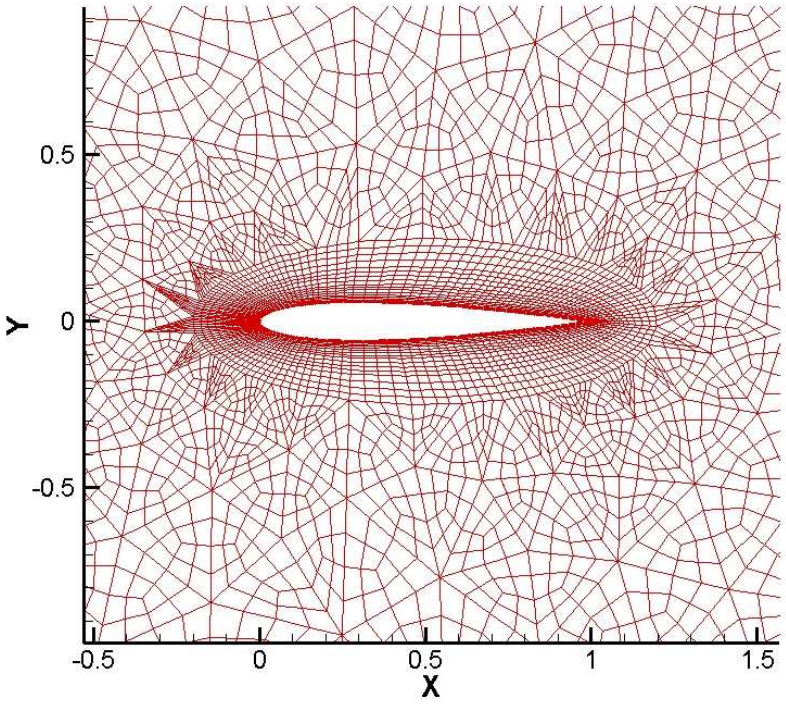


(b) the 4th order SD method;

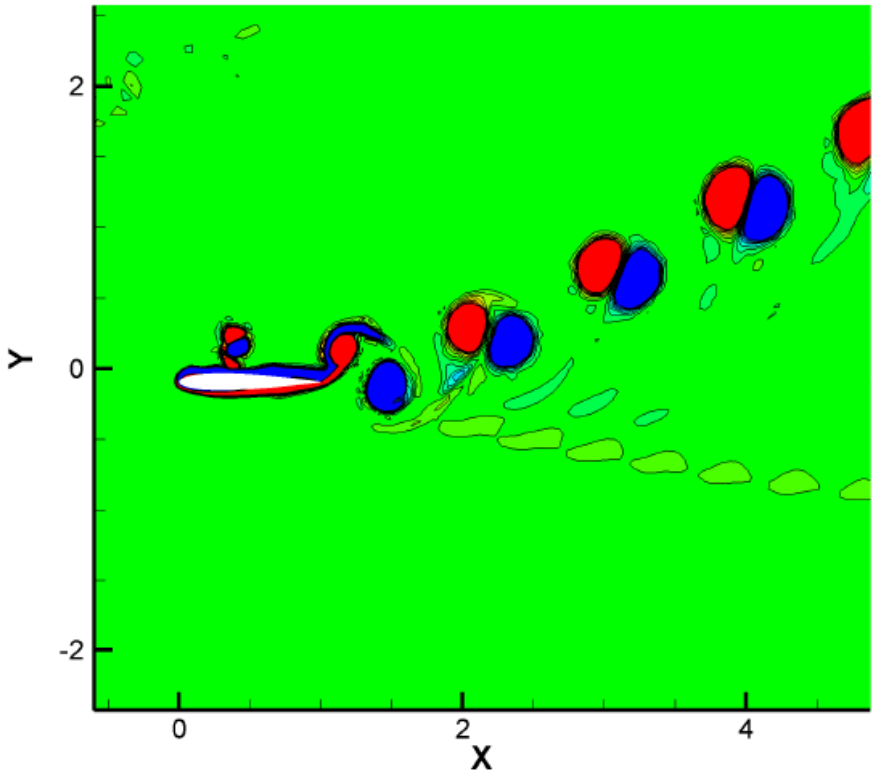
Mean drag = -0.51

Mean lift = 2.57

Vorticity shedding on the finest mesh

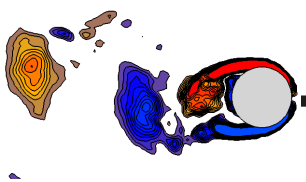


(a) Mesh C



(b) Vorticity $-6 \leq \omega c/U_\infty \leq 6$

Third-order SD method is used



Conclusions

- The SD method is robust and able to model compressible viscous flow with moving boundaries with high accuracy.
- It enables accurate simulations of vortex dominated flows
- Potential applications include flapping wing flight and wind turbine aerodynamics